

Probing quantum matter via randomized measurements



Wikipedia

B. Vermersch (University of Innsbruck)

with A. Elben, M. Dalmonte, J. I Cirac, T. Brydges, C. Maier, P. Jurcevic, N. Lanyon, R. Blatt, L. Sieberer, J. Yu, G. Zhu, N. Yao, M. Hafezi, and P. Zoller



UNIVERSITY OF INNSBRUCK



OAW

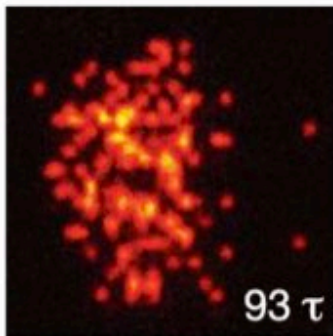
UQUAM
ERC Synergy Grant



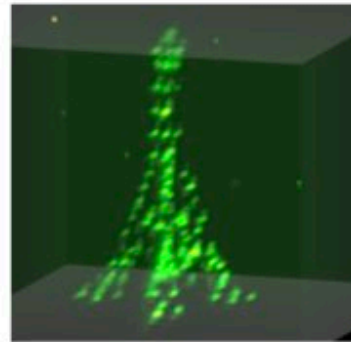
Research at Innsbruck...



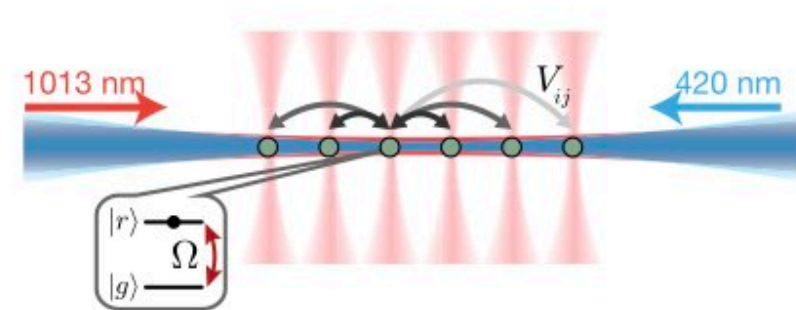
Ultracold atoms – Rydberg atoms



Choi et al., Science (2016)

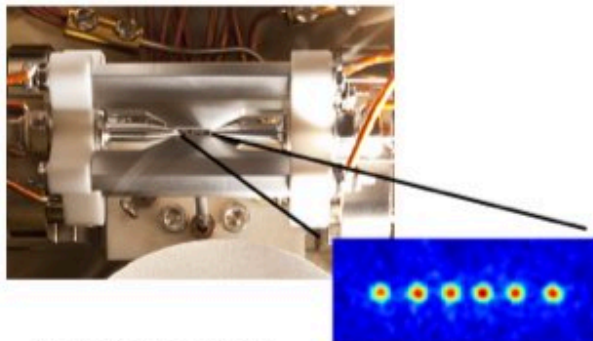


Barredo et al., Science (2016)



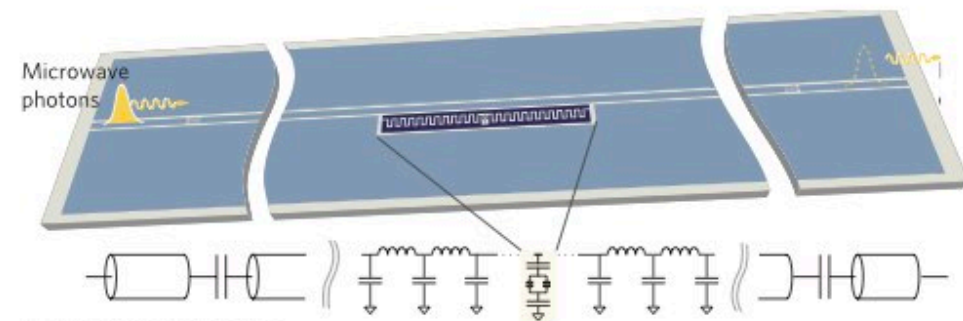
Bernien et al Nature 551, 579 (2017).

Trapped Ions



R. Blatt, Innsbruck

Superconducting circuits



A. Houck, Princeton

and quantum dots, NV centers, cavity QED, ..

Unique ways to create, probe, and understand quantum matter

Introducing Random measurements: Measuring Entanglement

Elben, Vermersch et al. (PRL 2018)
 Vermersch, Elben et al. (PRA 2018)
 arXiv:1806.05747
 arXiv:1812.02624



Measuring scrambling

arxiv: 1807.09087



Classification of interacting topological phases

in preparation

SYMMETRY-BREAKING DOES NOT OCCUR			SYMMETRY-BREAKING OCCURS
Trivial	Odd-spin Haldane phase	Topological insulator	$Z \rightarrow ZZ$ (Dimer phase)
CZX halozetype state	G-symmetry-protected topological state		$Z2 \rightarrow 1$ (Ising ferromagnet phase)
			Additional phases not listed

Renyi entropies as an entanglement measure



Two subsystems A and B are **bipartite entangled** iff

$$|\Psi\rangle \neq |\Psi_A\rangle \otimes |\Psi_B\rangle$$

Reduced density matrix

$$\rho_A = \text{Tr}_B |\Psi\rangle \langle \Psi|$$

Purity of subsystems

$$\text{Tr} [\rho_A^2], \text{Tr} [\rho_B^2] < \text{Tr} [\rho^2]$$

Purity of full system

Entanglement entropies

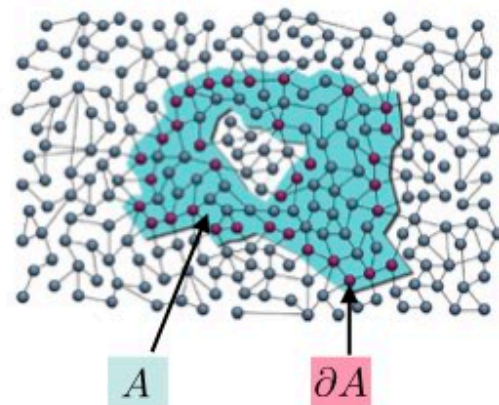
$$S_A = -\text{Tr}_A [\rho_A \log \rho_A]$$

von-Neumann

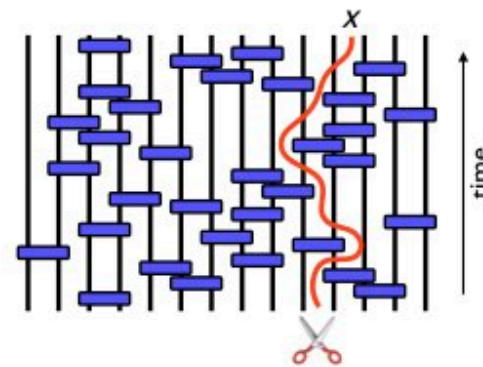
$$S_A^{(n)} = \frac{1}{1-n} \log \text{Tr}_A [\rho_A^n] \leq S_A$$

Renyi

Probe of Quantum Phases: Equilibrium,
Quantum phase transitions, Thermalization...



Eisert et al., Rev. Mod. Phys. 82, 277 (2010)



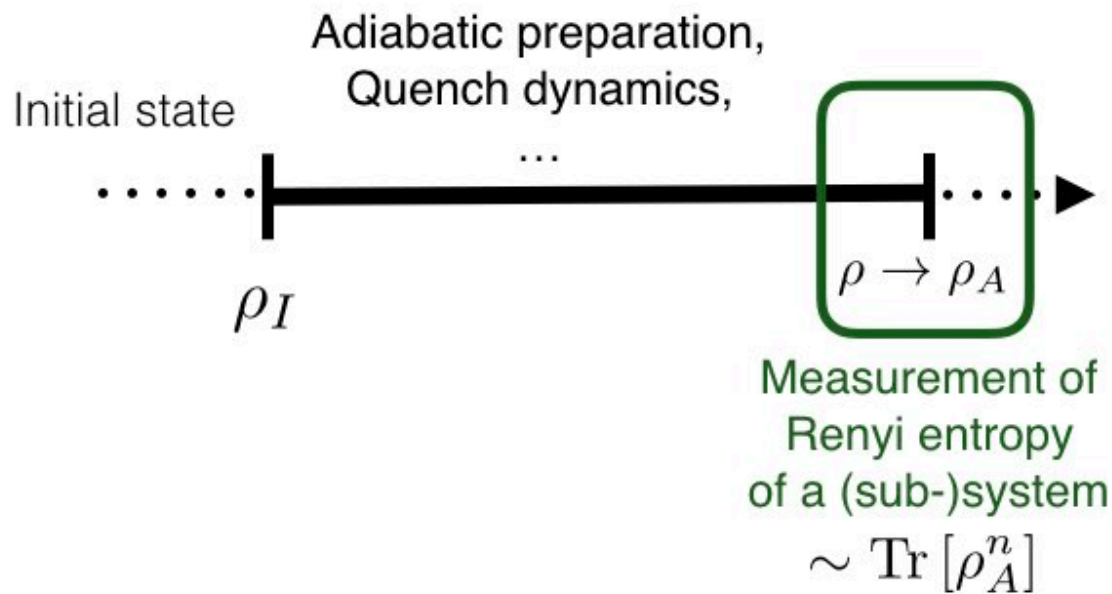
Nahum et al, Phys. Rev. X 7, 031016 (2017)

Verification of quantum simulators and quantum computers

Coherence, entanglement

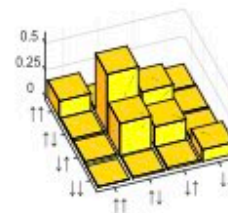
How to measure entanglement in (large) quantum systems?

How to measure Renyi entropies?



Realization in a trapped ion quantum simulator existing today in the lab?!

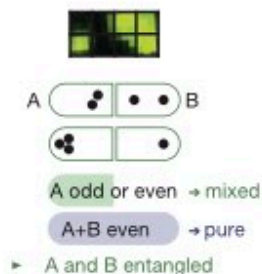
Quantum State Tomography



Gross et al. PRL 105, 150401 (2010)
B. P. Lanyon et al., Nat.Phys., (2017)

If available + feasible

Interference of copies

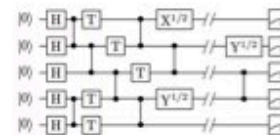


Islam et al., Nature 528, 77–83 (2015)

Daley et al. PRL, 109(2), 20505 (2012)

If available

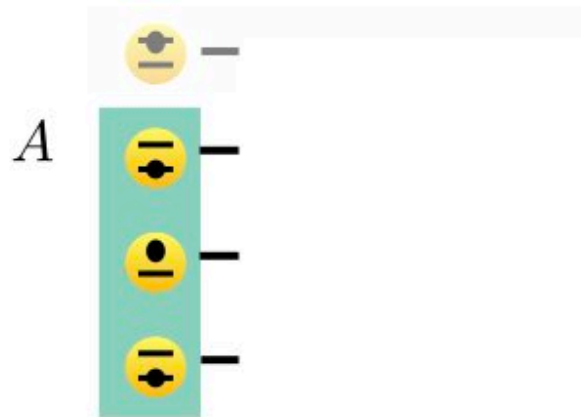
Random measurements on single copies in a QC



van Enk, Beenaker, PRL 108, 110503 (2012)

Random measurements in a *quantum computer*

Random measurement



$$P_U(\mathbf{s}_A) = \text{Tr} \left[U_A \rho_A U_A^\dagger |\mathbf{s}_A\rangle \langle \mathbf{s}_A| \right]$$

with $U_A \in \text{CUE}$

Average over the circular unitary ensemble (CUE)

$$\overline{P_U(\mathbf{s}_A)^2} = \frac{1 + \text{Tr} [\rho_A^2]}{N_{\mathcal{H}_A} (N_{\mathcal{H}_A} + 1)}$$

Hilbertspace dimension of A \nearrow

CUE (2-design): $\frac{\overbrace{U_{ik} U_{il}^*} \overbrace{U_{im} U_{in}^*}}{\sim \text{Gaussian}} = \frac{\delta_{kl} \delta_{mn} + \delta_{kn} \delta_{ml}}{N_{\mathcal{H}_A} (N_{\mathcal{H}_A} + 1)}$

Single qubit

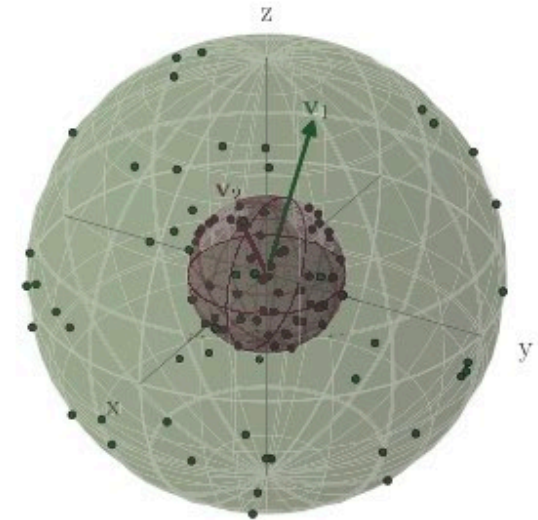
Bloch vector representation

$$\rho = \frac{1}{2} (\mathbb{1}_2 + \vec{a} \cdot \vec{\sigma})$$

$$\longrightarrow \text{Tr} [\rho^2] = \frac{1}{2} (1 + |\vec{a}|^2)$$

■ : $|\vec{a}| = 1$

■ : $|\vec{a}| = 0.33$



Random Measurement

$$\text{Rotation: } U\rho U^\dagger = \frac{1}{2} (\mathbb{1}_2 + (Q_U \vec{a}) \cdot \vec{\sigma})$$

$$\text{z-Measurements: } Z_U = \langle \sigma_z \rangle_U = (Q_U \vec{a})_3$$

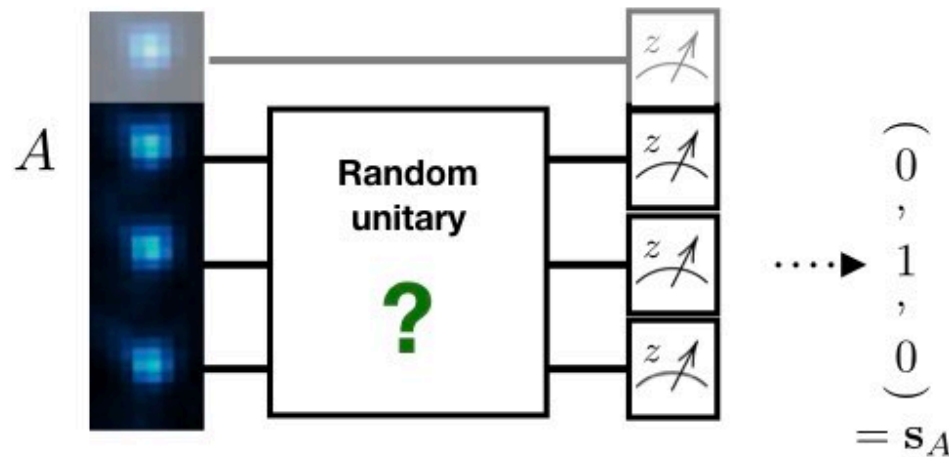
Z_U uniformly distributed with variance $\frac{|\vec{a}|^2}{3}$

$$\longrightarrow \text{Tr} [\rho^2] = \frac{1}{2} \left(1 + 3 \overline{\langle \sigma_z \rangle_U^2} \right)$$

The signal is in the noise

Random measurements in a *quantum simulator*

Random measurement in an experiment



$$P_U(\mathbf{s}_A) = \text{Tr} \left[U_A \rho_A U_A^\dagger |\mathbf{s}_A\rangle \langle \mathbf{s}_A| \right]$$

with $U_A \in \text{CUE}$

Average over the circular unitary ensemble (CUE)

$$\overline{P_U(\mathbf{s}_A)^2} = \frac{1 + \text{Tr} [\rho_A^2]}{N_{\mathcal{H}_A} (N_{\mathcal{H}_A} + 1)}$$

Hilbertspace dimension of A \nearrow

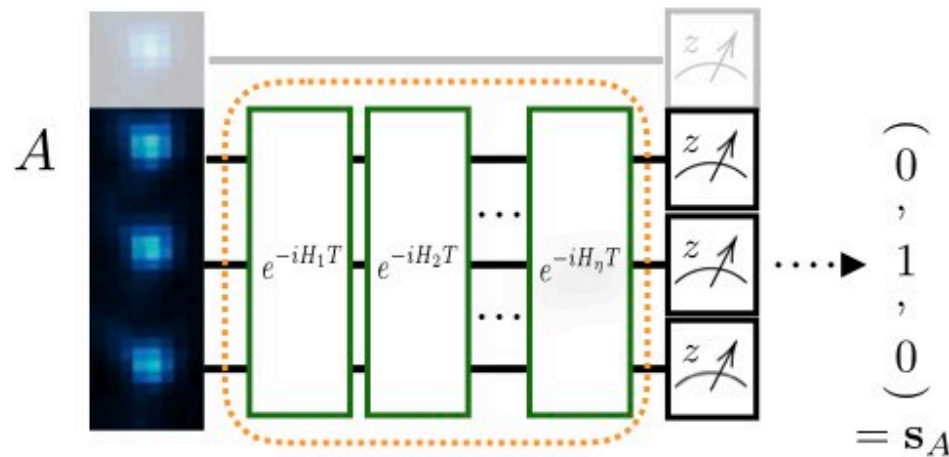
How to generate **random unitaries from CUE?**

A **simpler** protocol?

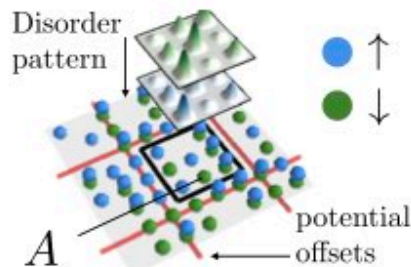
How many **measurements** per unitary and **how many unitaries?**

Random measurements via random quenches

Random unitaries via random quenches



Hamiltonian time evolution with
✓ generic interactions
✓ engineered, *timedependent* disorder



$$P_U(\mathbf{s}_A) = \text{Tr} \left[U_A \rho_A U_A^\dagger |\mathbf{s}_A\rangle \langle \mathbf{s}_A| \right]$$

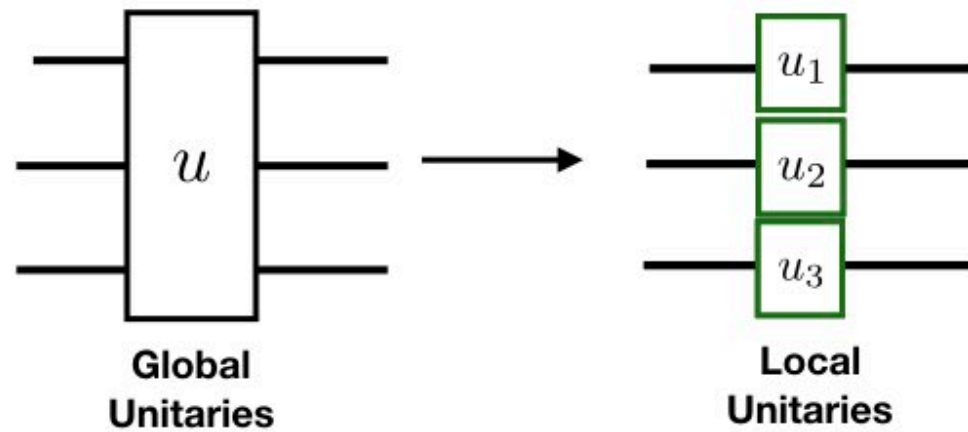
with

$$U_A = e^{-iH_n T} \dots e^{-iH_1 T}$$

**Applicable in Bose/Fermi
Hubbard and Spin models**

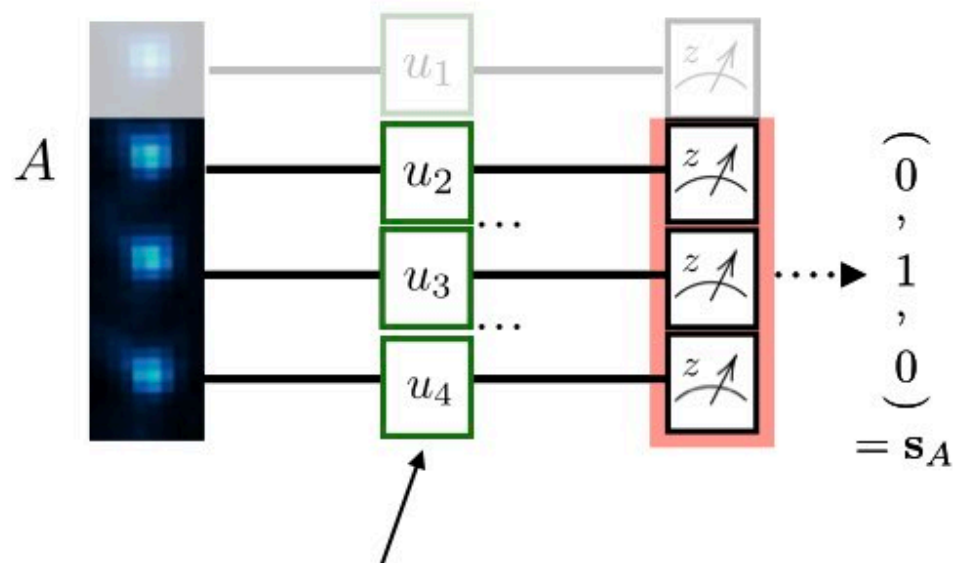
- Ultracold atoms
- Rydberg atoms
- Superconducting qubits
- Trapped ions
- ...

A simpler protocol?

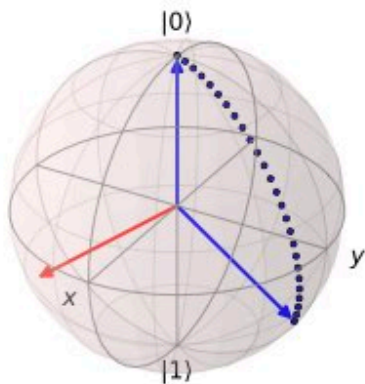


Random measurements with *local* unitaries

Random measurements with local unitaries



Single spin random rotation



$$P_U(\mathbf{s}_A) = \text{Tr} \left[U_A \rho_A U_A^\dagger |\mathbf{s}_A\rangle \langle \mathbf{s}_A| \right]$$

$$U_A = \bigotimes u_i \quad u_i \in \text{CUE}(d)$$

$$u_i \equiv u(\alpha_i, \beta_i, \gamma_i)$$

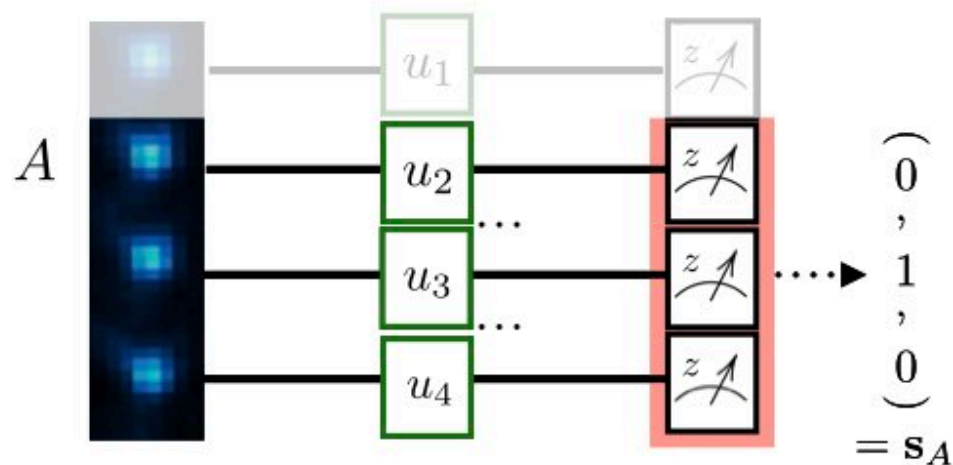
$$= Z(\alpha_i) Y(\pi/2) Z(\beta_i) Y(-\pi/2) Z(\gamma_i)$$

$\alpha_i, \beta_i, \gamma_i$ random

Elben, Vermersch et al. (PRL 2018)
Brydges, Elben et al (arXiv:1806.05747)

Random measurements with *local* unitaries

Random measurements with local unitaries



$$P_U(\mathbf{s}_A) = \text{Tr} \left[U_A \rho_A U_A^\dagger |\mathbf{s}_A\rangle \langle \mathbf{s}_A| \right]$$

$$U_A = \bigotimes u_i \quad u_i \in \text{CUE}(d)$$

Average over the random unitary ensemble

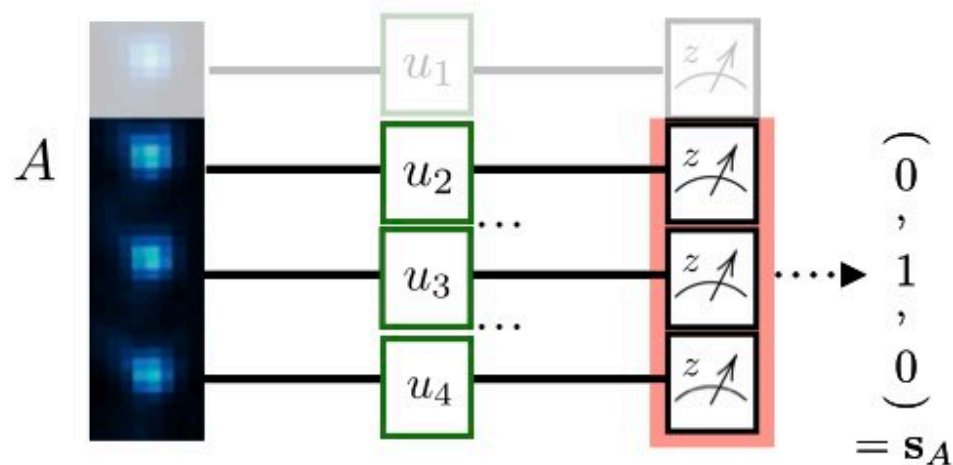
$$\overline{P_U(\mathbf{s}_A)} = \frac{1}{d^{N_A}} \quad \overline{P_U(\mathbf{s}_A)^2} = \frac{\sum_{A' \subseteq A} \text{Tr}(\rho_{A'}^2)}{d^{N_A} (d+1)^{N_A}}$$

↑
Exponential number of terms

How to invert, to obtain directly $\text{Tr}[\rho_A^2]$?

Random measurements with *local* unitaries

Random measurement with local random unitaries



$$P_U(\mathbf{s}_A) = \text{Tr} \left[U_A \rho_A U_A^\dagger |\mathbf{s}_A\rangle \langle \mathbf{s}_A| \right]$$

$$U_A = \bigotimes u_i \quad u_i \in \text{CUE}(d)$$

Average over the random unitary ensemble

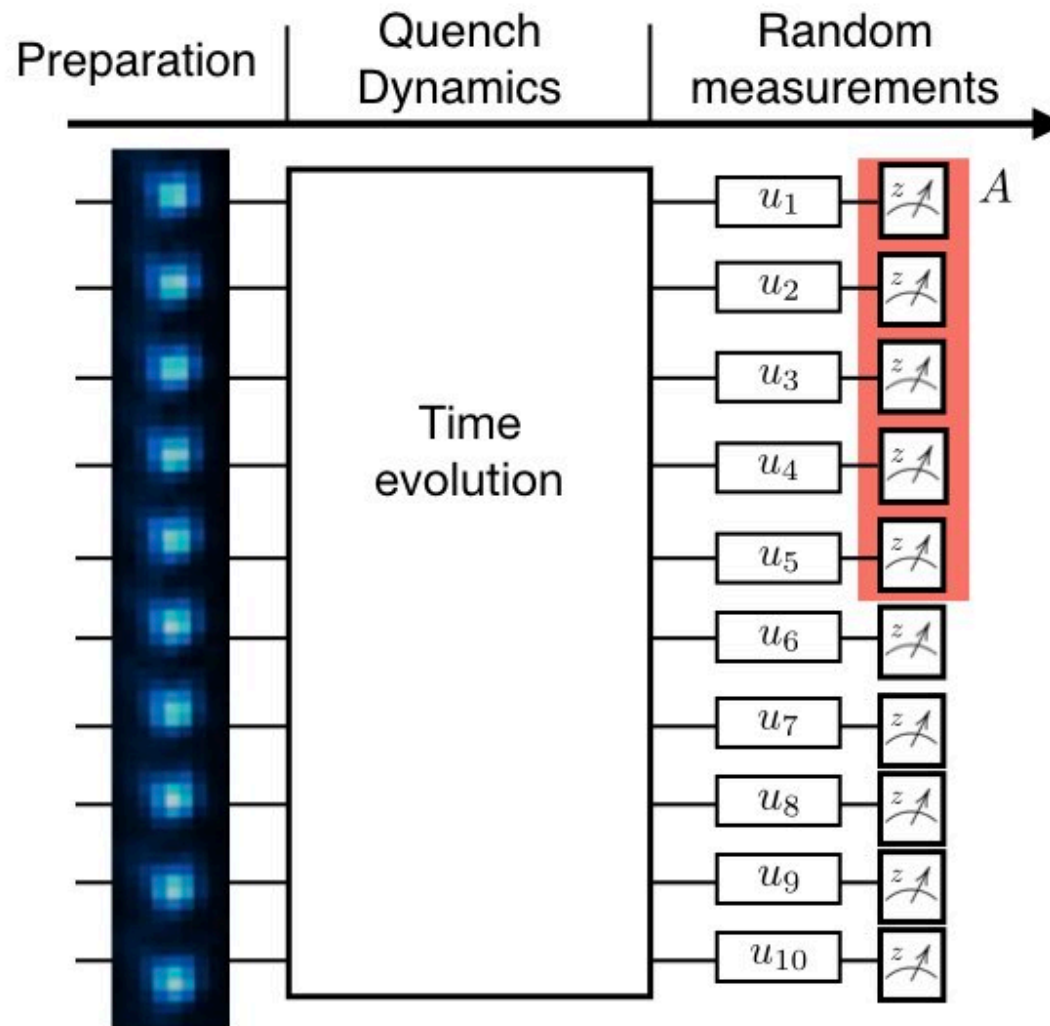
$$\text{Tr} [\rho_A^2] = \overline{X_U} \quad \text{with}$$

Explicit reconstruction

$$X_U = 2^{N_A} \sum_{s_A, s'_A} (-2)^{-D[s_A, s'_A]} \underbrace{P_U(s_A) P_U(s'_A)}_{\text{Cross correlation}}$$

Hamming distance

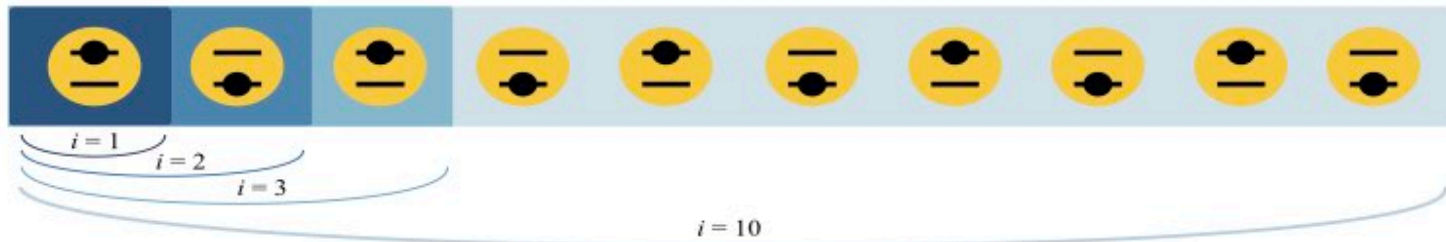
Trapped ion measurement scheme



$$H_{XY} = \hbar \sum_{i < j} J_{ij} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+) + \hbar \sum_j (B + b_j) \sigma_j^z$$

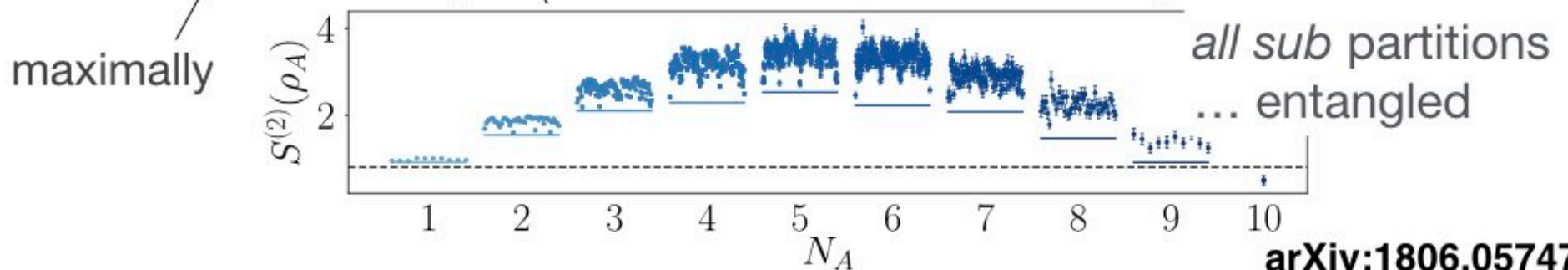
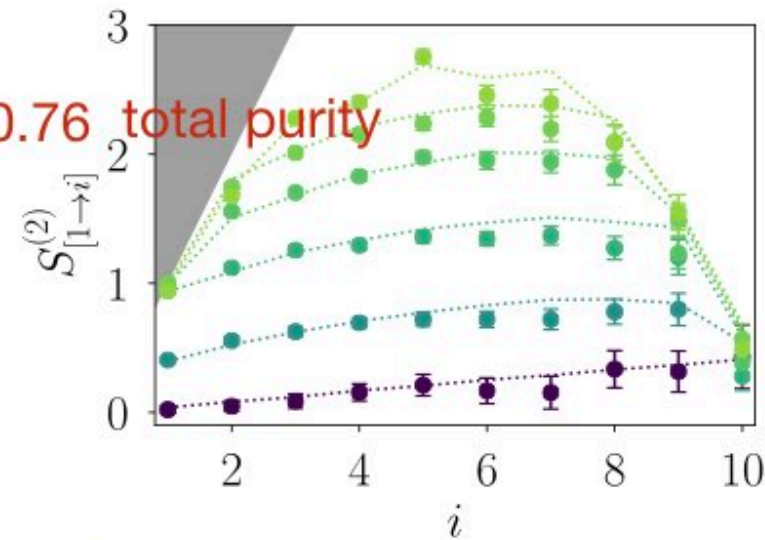
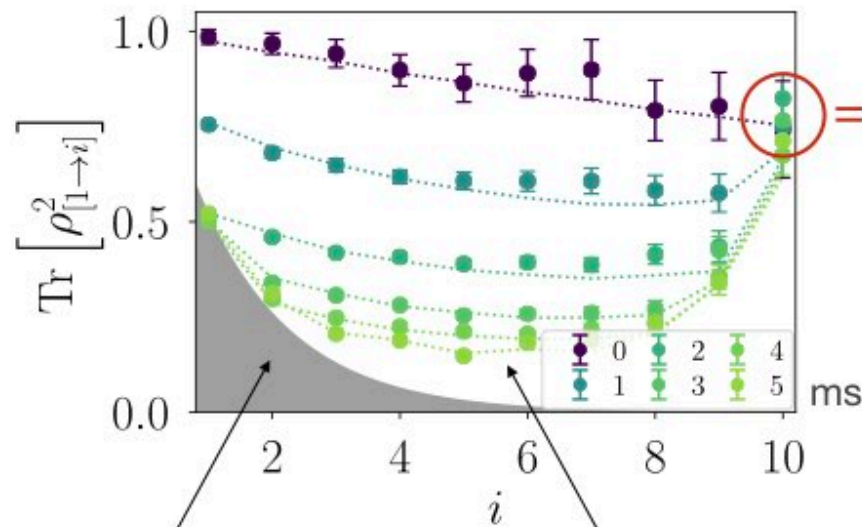
long range (algebraic) interaction
local disorder potentials

Experimental results - 10 ions [no disorder]



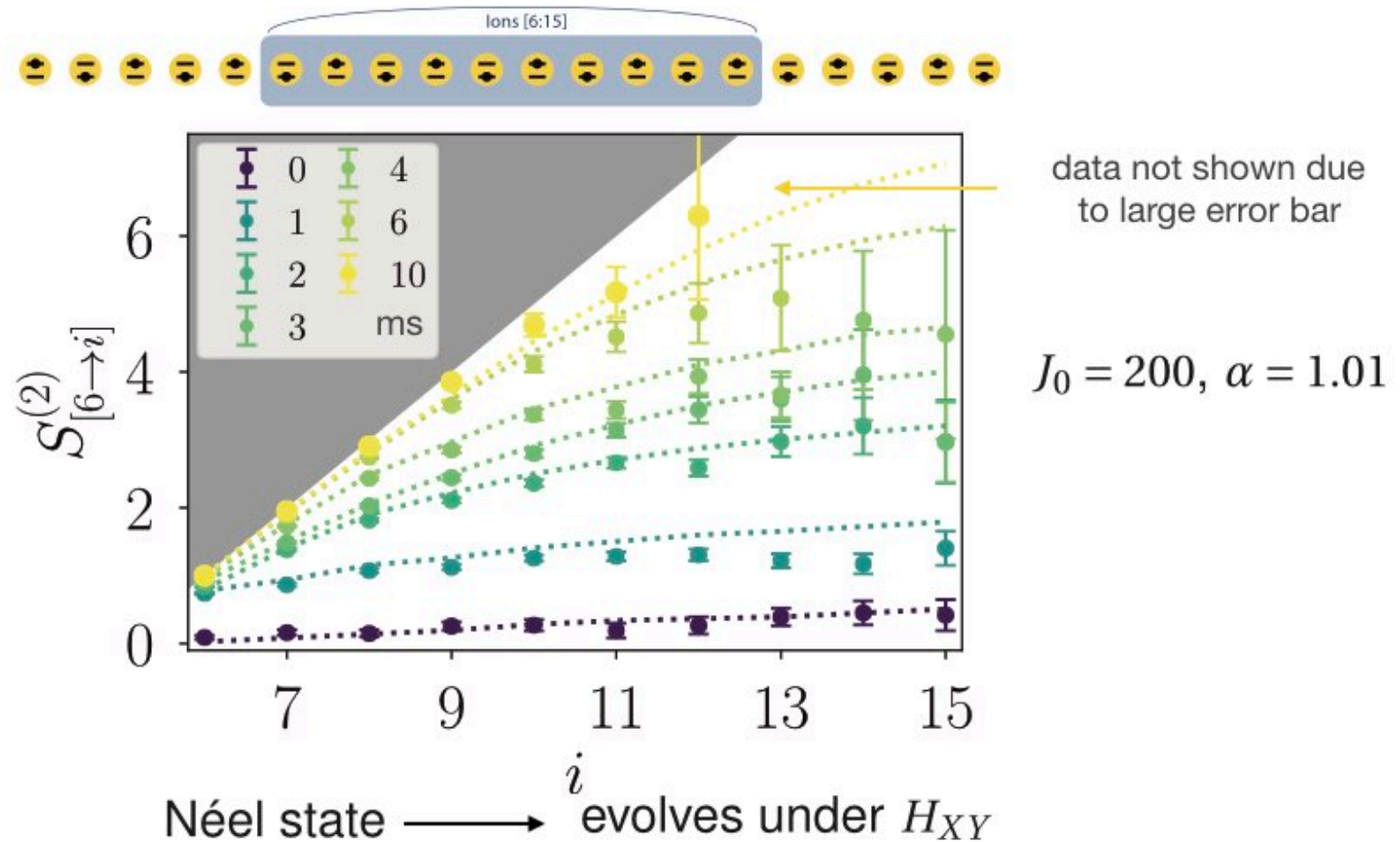
purity $\text{Tr}[\rho_A^2]$

Renyi entropy



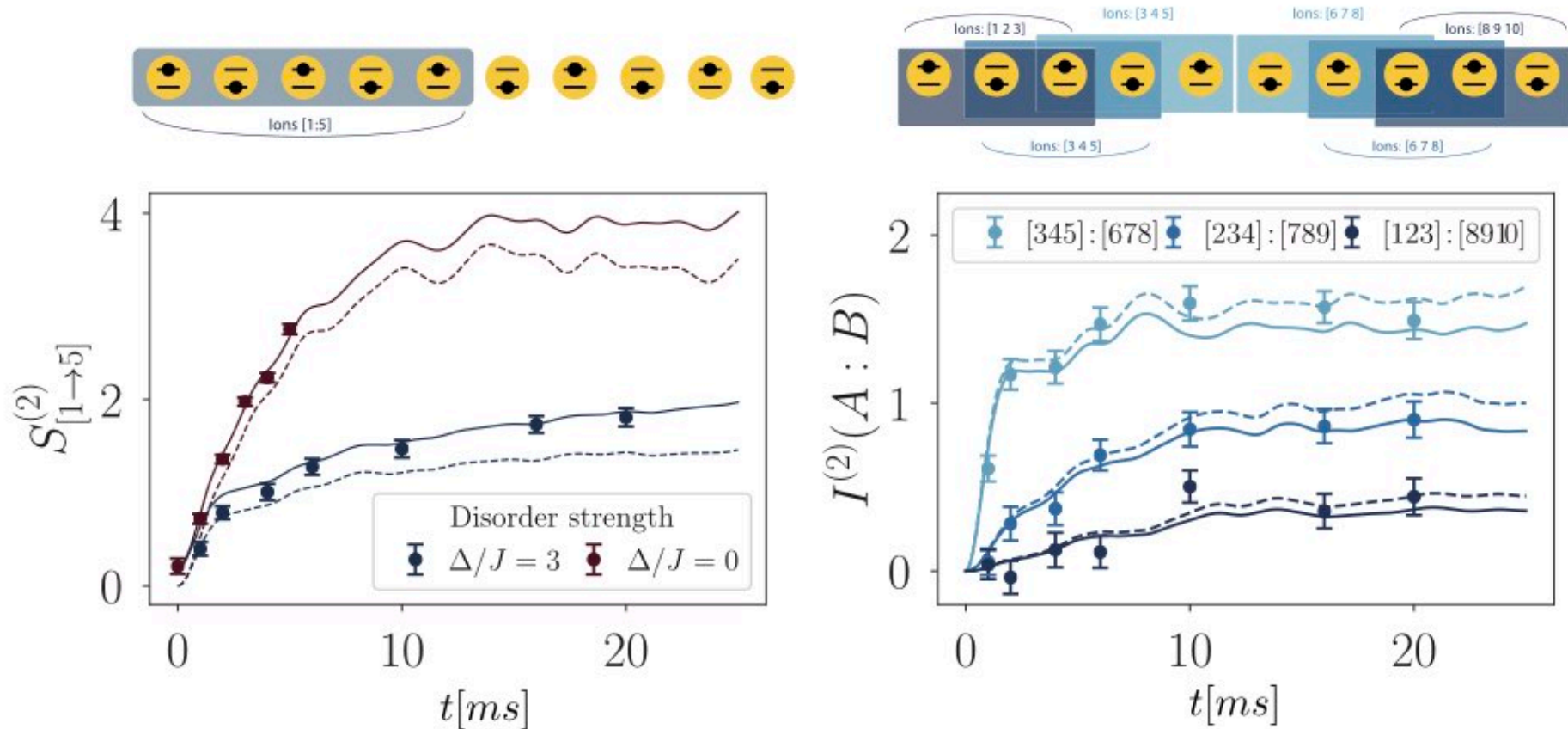
Experimental results - 20 Ions [no disorder]

Renyi entropy 1-10-qubit partitions of a 20-qubit system.



With minor improvements of the apparatus Renyi entropies of 20 ions seem accessible

Experimental results - 10 ions [disorder]

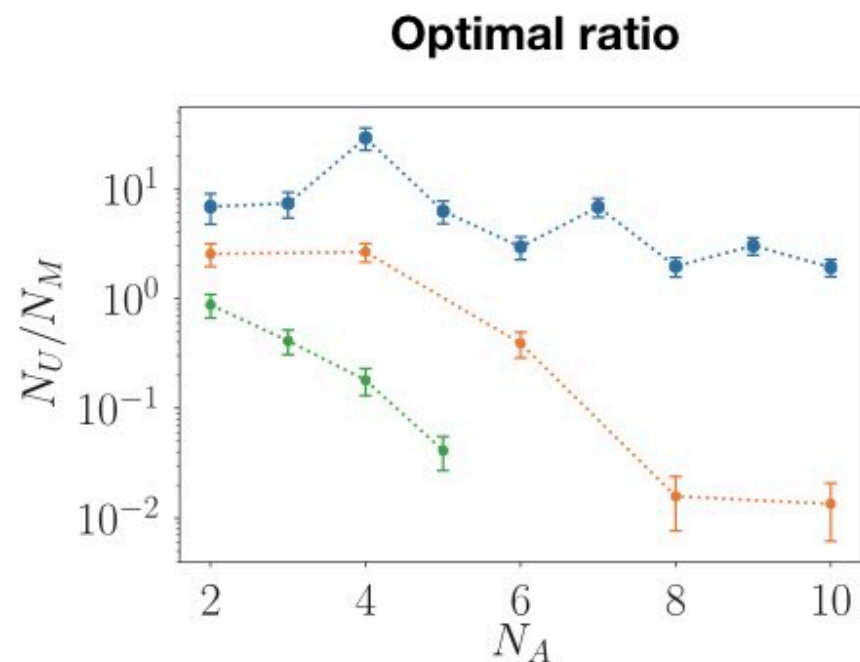
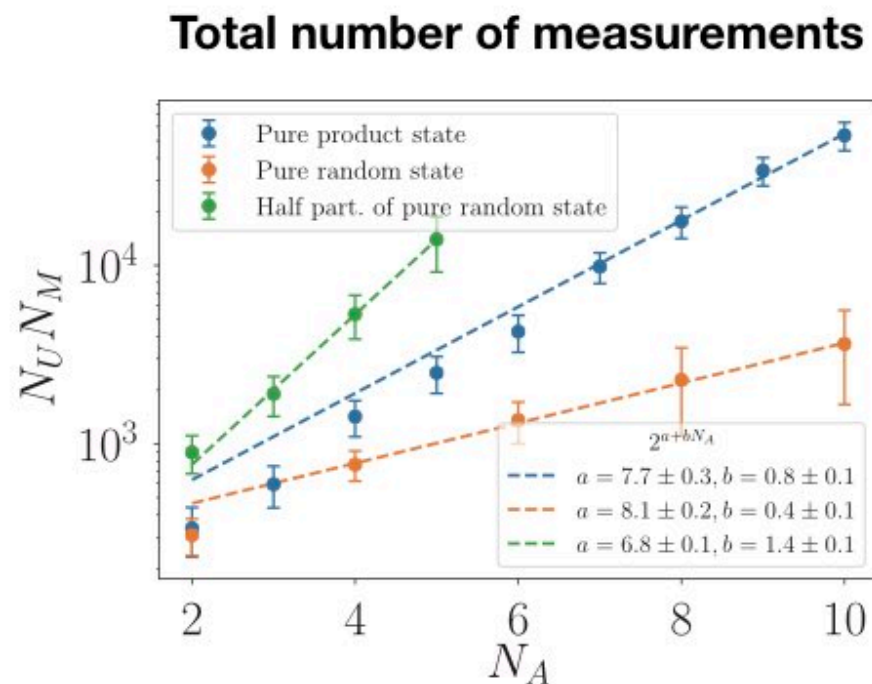


$$H_{XY} = \hbar \sum_{i < j} \boxed{J_{ij}} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+) + \hbar \sum_j (B + \boxed{b_j}) \sigma_j^z$$

long range interaction
local disorder potentials

Scaling of the required number of measurements

Scaling of the number of measurements to obtain the Renyi entropy with an error smaller than 0.15

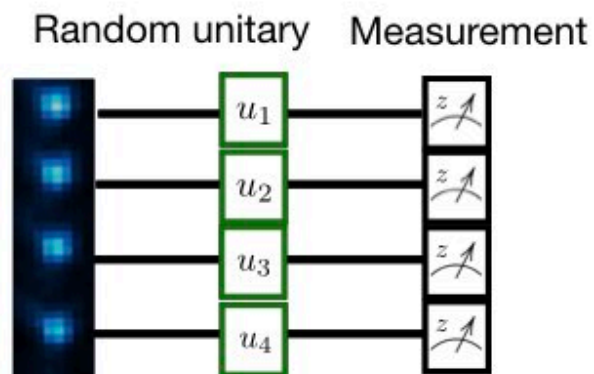


- Exponential scaling with exponents < 1.5
- Optimal ratio strongly dependent on the state of interest

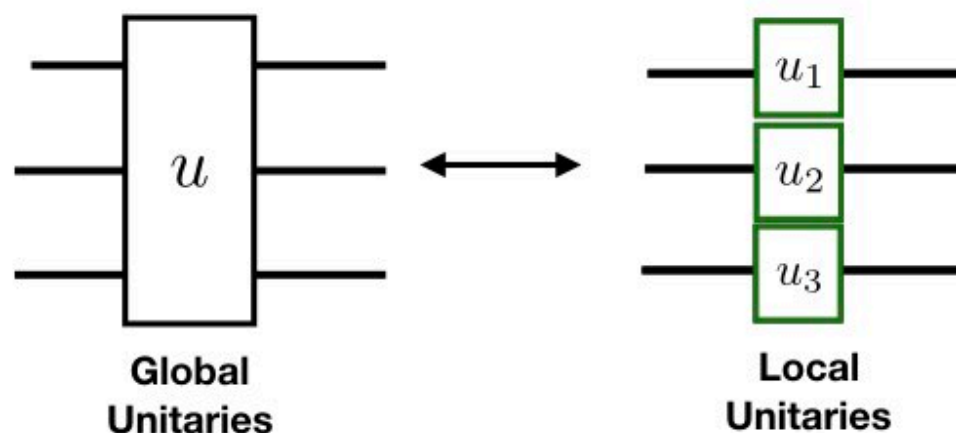
→ Adaptive sampling to find optimal parameters?

Our tool: random measurements

- Single copies
- Local Measurements
- Robust (ex: readout errors)
- Simple



- Two variants



A powerful tool to probe quantum matter?

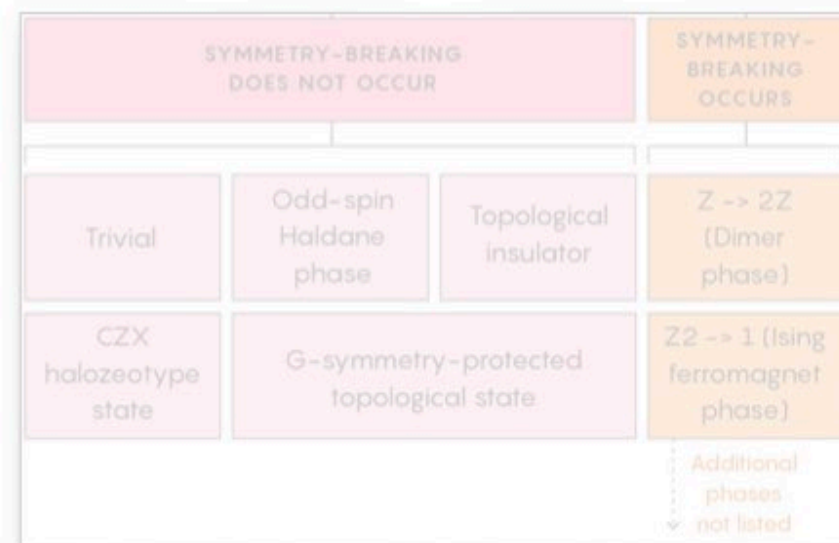
Measuring scrambling with random measurements



**B. Vermersch, A. Elben, L. Sieberer,
N. Yao, and P. Zoller**

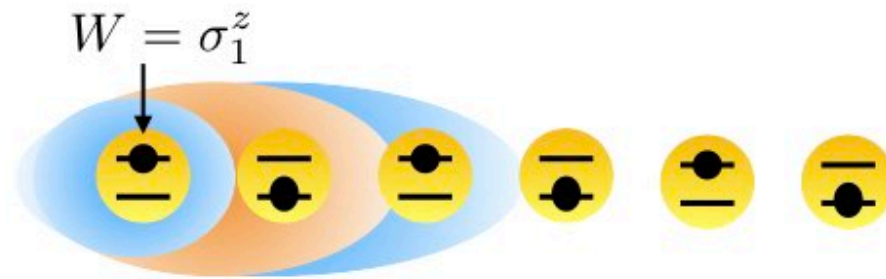
arxiv: 1807.09087

Classification of interacting topological phases (SPT)



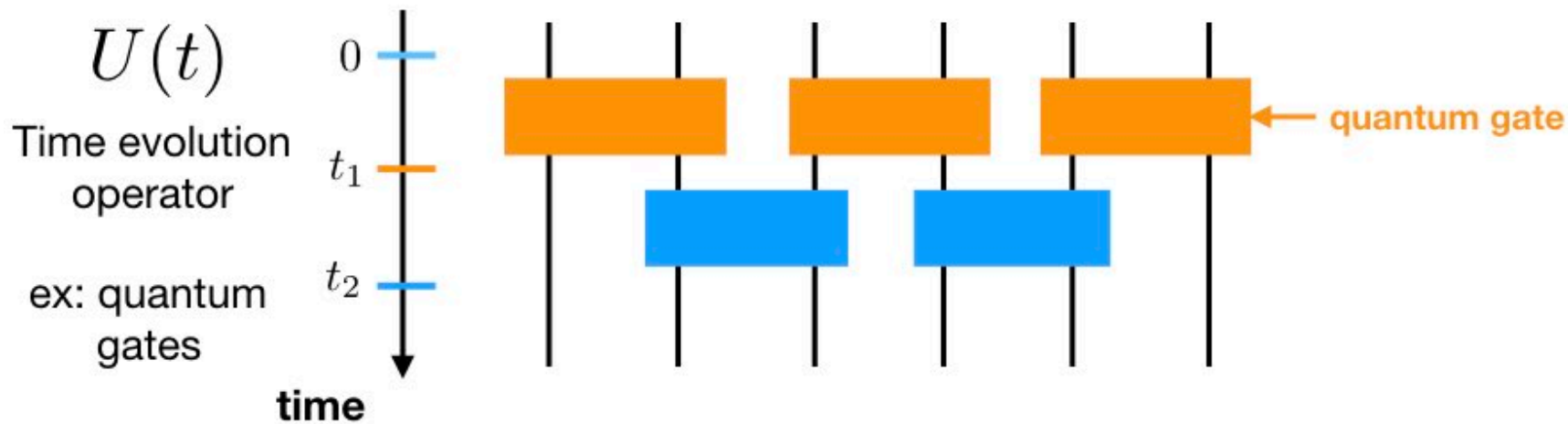
**A. Elben, B. Vermersch, J. Yu, G. Zhu,
M. Hafezi and P. Zoller**

in preparation



Operator spreading
(Heisenberg picture)

$$W(t) = U^\dagger(t) W U(t)$$



$$W(0) = \sigma_1^z \otimes I \otimes \dots$$

$$W(t_1) = \sum_{\gamma_1, \gamma_2 = x, y, z, 0} c_{\gamma_1, \gamma_2} \sigma_1^{\gamma_1} \otimes \sigma_2^{\gamma_2} \otimes I \dots$$

⋮

OTOC

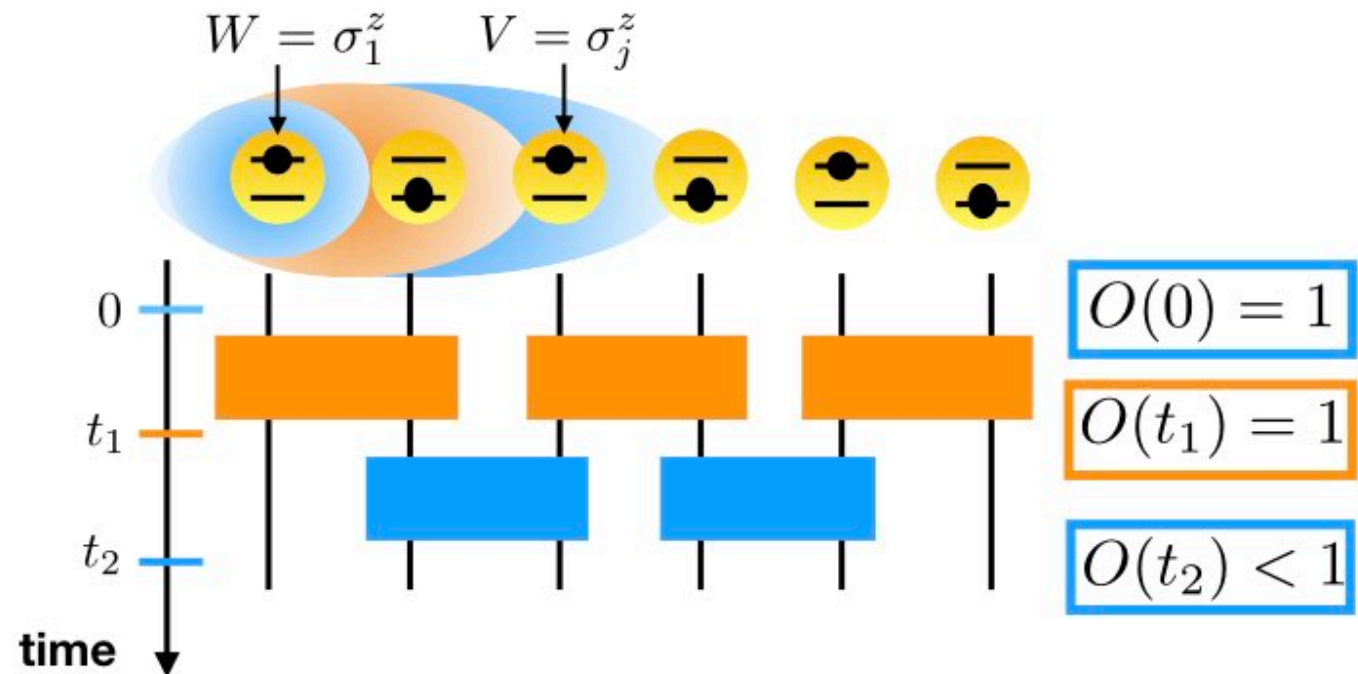
$$O = \text{Tr}(\rho W(t) V W(t) V)$$

↑
quantum state (here: V, W Pauli matrices)

$$[W(t), V] = 0 \rightarrow O = 1$$

$$[W(t), V] \neq 0 \rightarrow O < 1$$

Describe the spreading of an operator with respect to a 'reference' V



OTOCs: Information scrambling in many-body systems

Fast scramblers
(analog black holes)

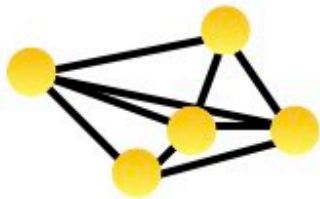
Quantum
Thermalization
And chaos

Slow scramblers

Sachdev-Ye-Kitaev (SYK) model

Random circuits

Many-body localization

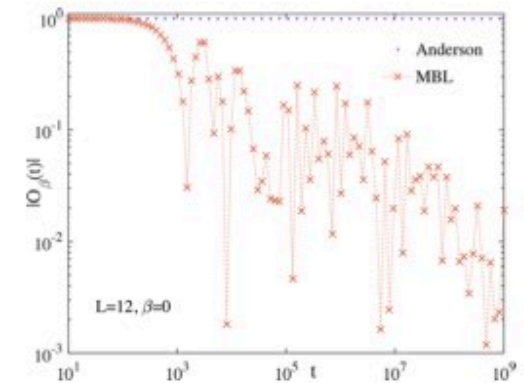
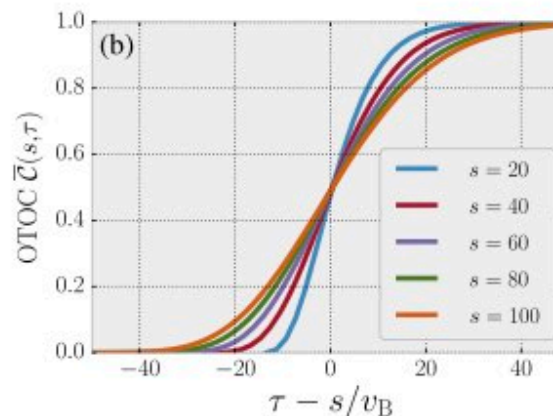


Quantum Ising models
(inc. Long rang interactions)
Bose Hubbard

$$H_{SYK} = \frac{1}{4!} \sum_{jklm} J_{jklm} \chi_j \chi_k \chi_l \chi_m$$

$$O(t, r) \sim c_0 - c_1 e^{\lambda_L(t-r/v_B)}$$

$$\lambda_L = 2\pi k_B T / \hbar$$



Sachdev, et al . PRL 1993 70(21), 3339–3342
Kitaev, A. KITP 2015
Banerjee et al 2017 PRB 95(13), 134302.

A. Bohrdt et al, New J. Phys. 19, 063001 (2017).
A. Nahum et al Phys. Rev. X 8, 021014 (2018).
C. W. von Keyserlingk et al Phys. Rev. X 8, 021013 (2018).
M. C. Tran, et al A. V. Gorshkov, arxiv:1808.05225 .

Fan, et al . Science Bulletin, 62(10), 707–711
Chen, X et al. Annalen Der Physik, 529(7)

...

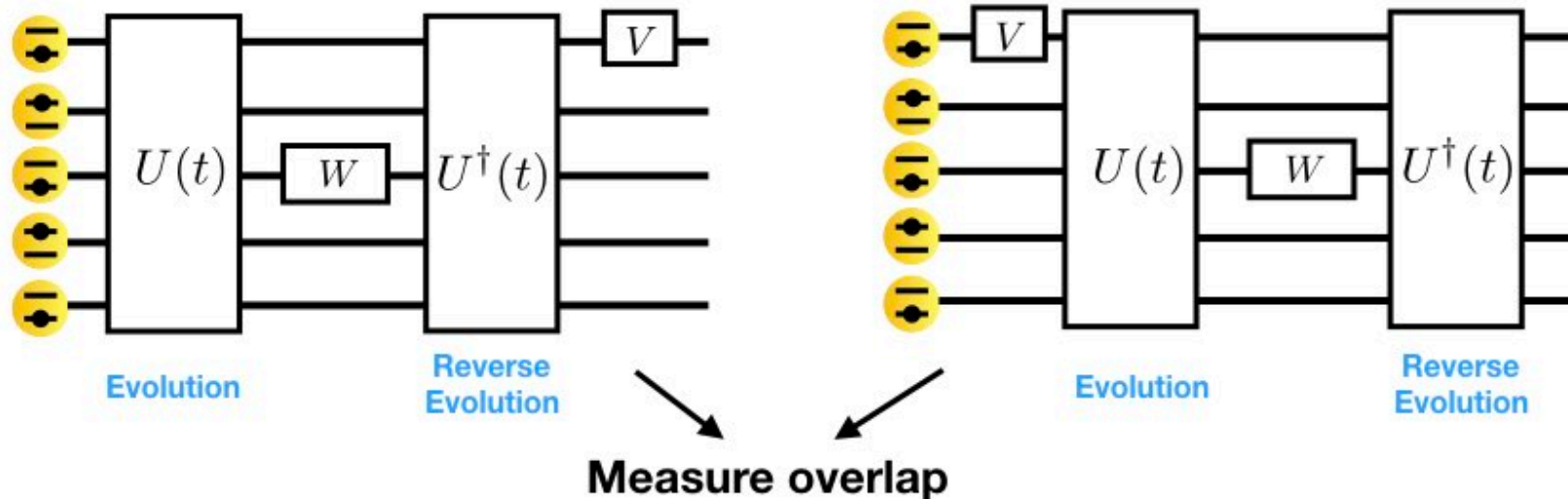
Peculiar time-ordering in the definition implies 'challenging' protocols

$$O = \text{Tr}(\rho W(t) V W(t) V)$$

For a pure state $\rho = |\psi\rangle\langle\psi|$ $O = \langle\psi_1|\psi_2\rangle$

$$|\psi_1\rangle = V U^\dagger(t) W U(t) |\psi\rangle$$

$$|\psi_2\rangle = U^\dagger(t) W U(t) V |\psi\rangle$$



Requires time-reversal and/or copies and/or ancillas

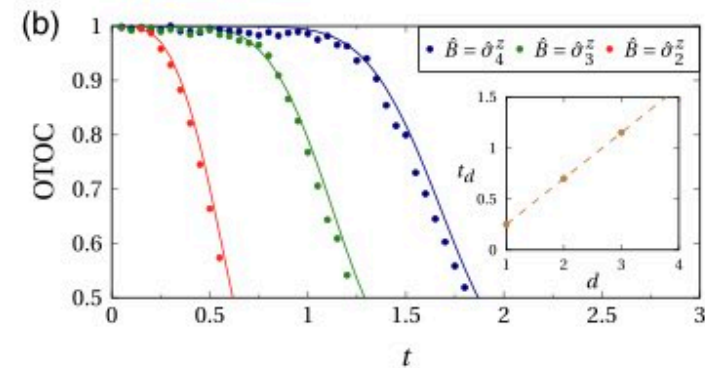
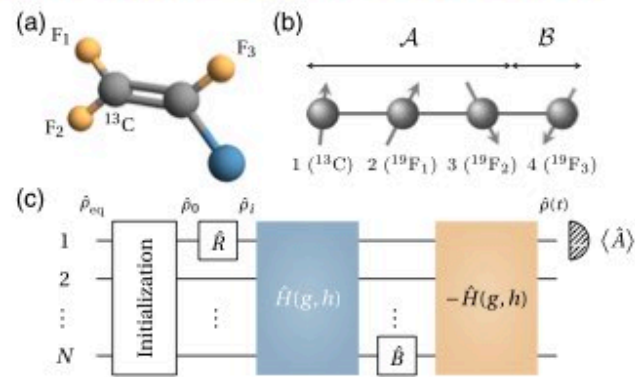
Zhu et al. *Phys. Rev. A* 94 062329 (2016)

Swingle, B. et al *Phys. Rev. A* 94, 1–6 (2016)

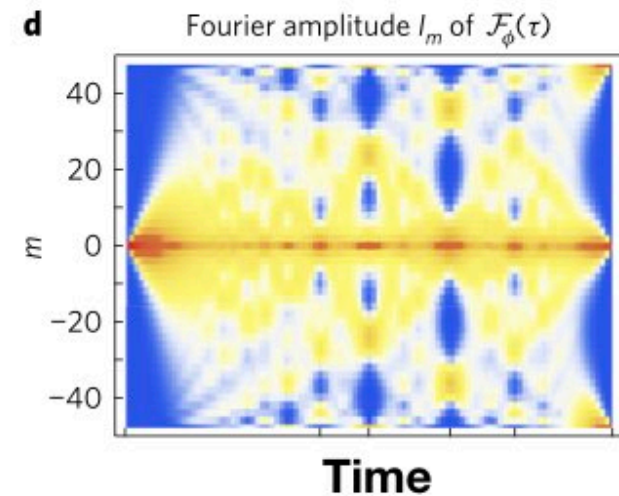
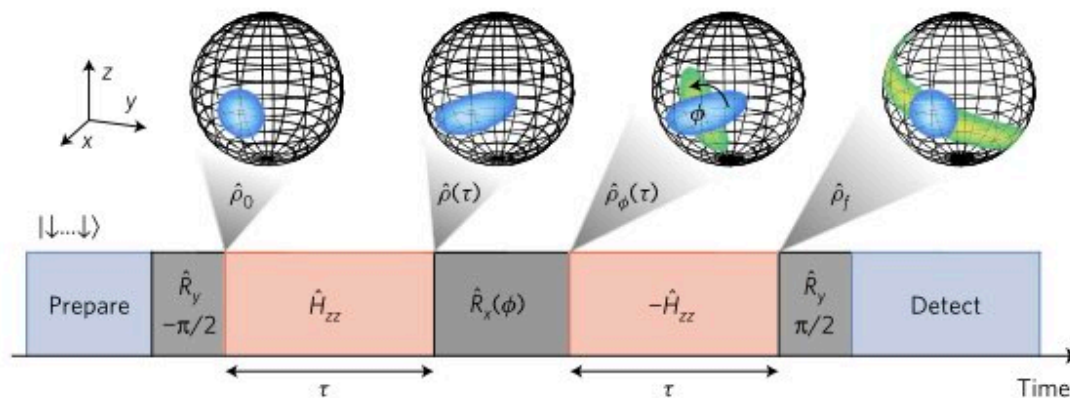
Yao et al, *arxiv:1607.01801*

Gartner, M., et al *Nature Physics*, 13(8), 781–786

NMR Four spins Trotter evolution

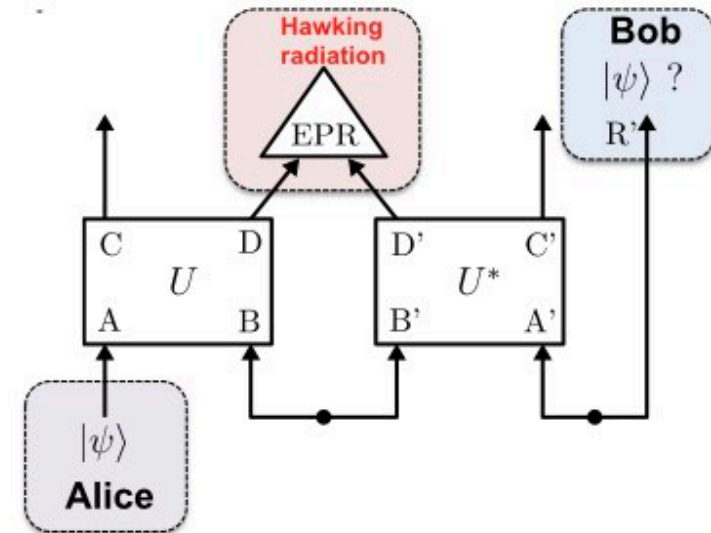
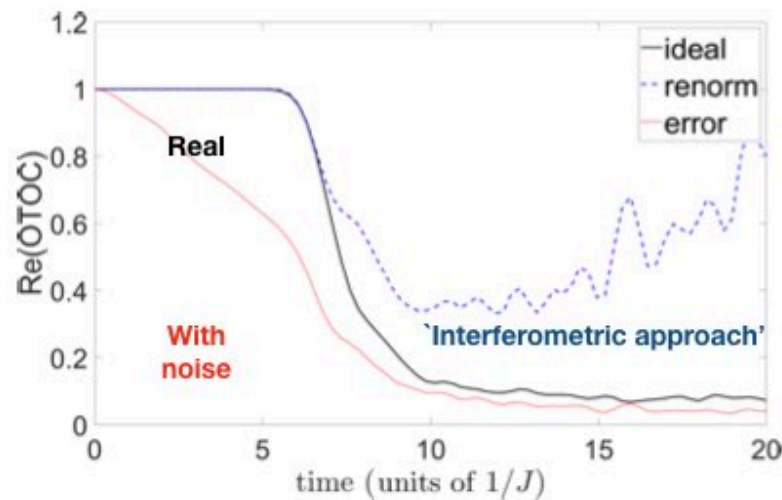


Trapped ions (all-to-all interactions)



Key challenges: \rightarrow Implementing time-reversal \rightarrow The role of decoherence

Decoherence versus scrambling



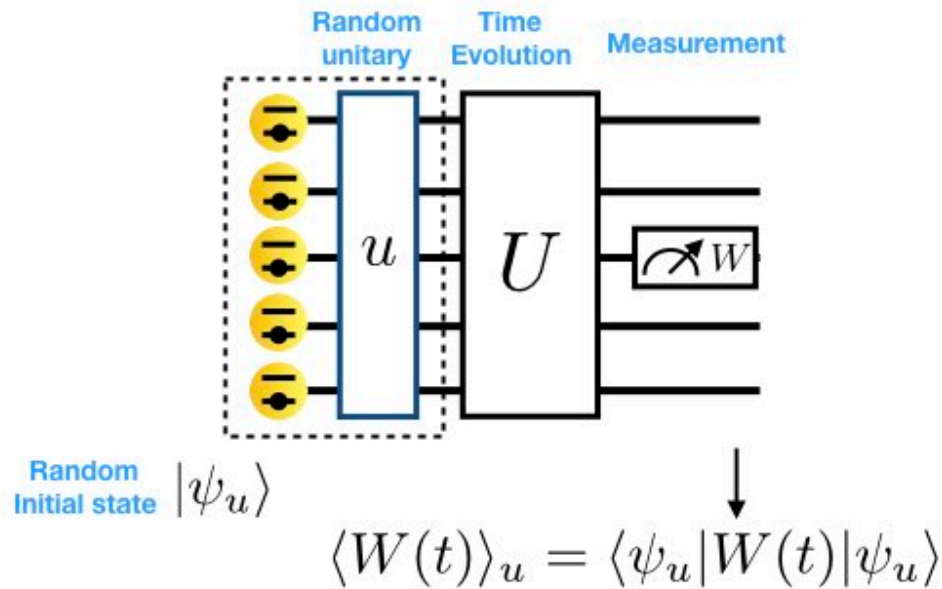
B. Swingle and N. Yunger Halpern, *Phys. Rev. A* 97, 062113 (2018).

B. Yoshida and N. Y. Yao, *arXiv:1803.10772*

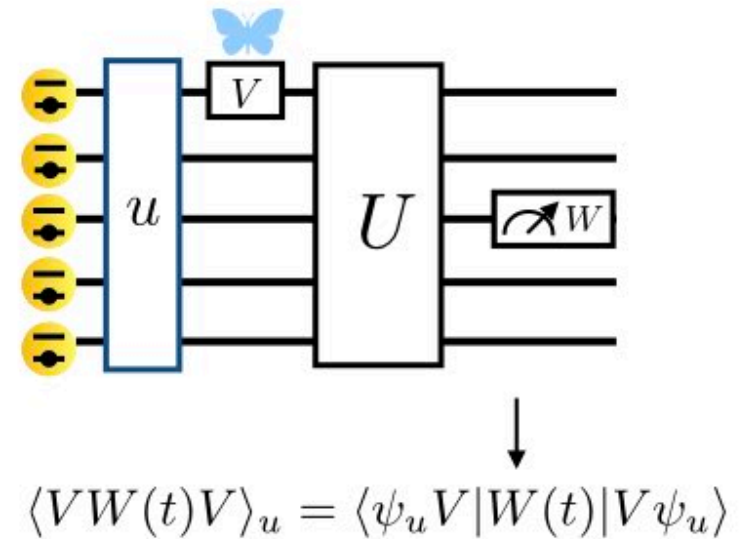
K. A. Landsman et al, *arxiv: 1806.02807*

- ➔ Very important technological **challenges..**
- ➔ **Our approach:** Replace time-reversal by **statistical correlations**

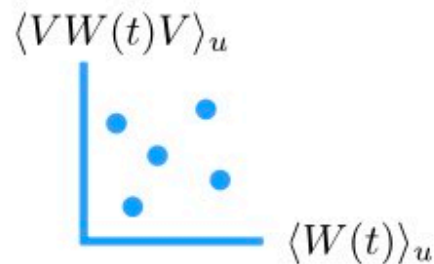
1st measurement (day 1)



2nd measurement (day 2)



Statistical correlations = OTOCs ($T = \infty$)



$$O(t) = \frac{1}{\mathcal{D}(G)} \overline{\langle W(t) \rangle_{u, k_0} \langle V^\dagger W(t) V \rangle_{u, k_0}}$$

ensemble average over the circular unitary ensemble (CUE)

CUE (2-design):

$$\begin{aligned}
 & \overline{u_{m_1, n_1} u_{m'_1, n'_1}^* u_{m_2, n_2} u_{m'_2, n'_2}^*} \quad (3) \\
 &= \frac{\delta_{m_1, m'_1} \delta_{m_2, m'_2} \delta_{n_1, n'_1} \delta_{n_2, n'_2} + \delta_{m_1, m'_2} \delta_{m_2, m'_1} \delta_{n_1, n'_2} \delta_{n_2, n'_1}}{\mathcal{N}_{\mathcal{H}}^2 - 1} \\
 &- \frac{\delta_{m_1, m'_1} \delta_{m_2, m'_2} \delta_{n_1, n'_2} \delta_{n_2, n'_1} + \delta_{m_1, m'_2} \delta_{m_2, m'_1} \delta_{n_1, n'_1} \delta_{n_2, n'_2}}{\mathcal{N}_{\mathcal{H}}(\mathcal{N}_{\mathcal{H}}^2 - 1)},
 \end{aligned}$$

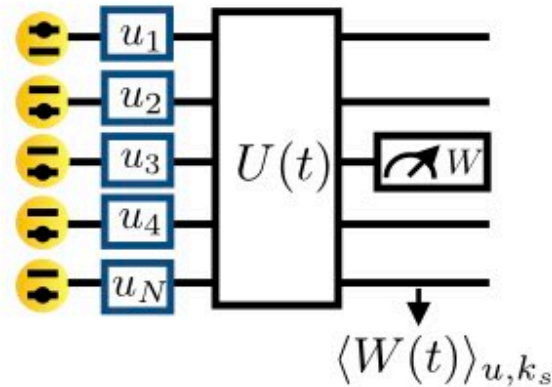
(Collins 2006)

$$\begin{aligned}
 \overline{\langle A \rangle_u \langle B \rangle_u} &= \overline{[u - \rho_0 - u^\dagger - A] [u - \rho_0 - u^\dagger - B]} \\
 &= \frac{1}{\mathcal{N}_{\mathcal{H}}^2 - 1} \left[\begin{array}{c} \overline{[u - \rho_0 - u^\dagger - A] [u - \rho_0 - u^\dagger - B]} \\ + \overline{[u - \rho_0 - u^\dagger - A] [u - \rho_0 - u^\dagger - B]} \end{array} \right] \quad \left. \begin{array}{l} \text{2-design} \\ \text{rule} \end{array} \right\} \\
 + \frac{-1}{\mathcal{N}_{\mathcal{H}}(\mathcal{N}_{\mathcal{H}}^2 - 1)} \left[\begin{array}{c} \overline{[u - \rho_0 - u^\dagger - A] [u - \rho_0 - u^\dagger - B]} \\ + \overline{[u - \rho_0 - u^\dagger - A] [u - \rho_0 - u^\dagger - B]} \end{array} \right] \\
 &= c \sum_{\tau \in I, \text{Swap}} \overline{\tau(A \otimes B)} = c \sum_{\tau \in I, \text{Swap}} \text{Tr}[\tau(A \otimes B)] = c \text{Tr}(AB)
 \end{aligned}$$

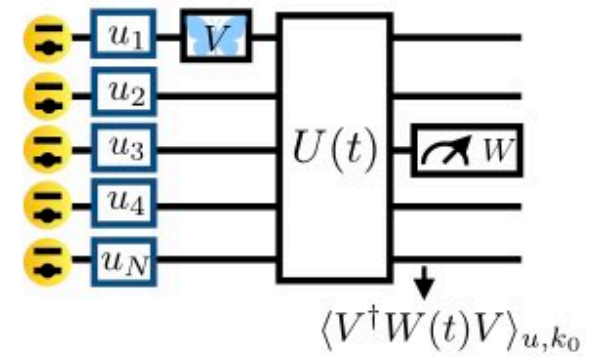
$$O(t) = \frac{1}{\mathcal{D}(\mathcal{G})} \overline{\langle W(t) \rangle_{u, k_0} \langle V^\dagger W(t) V \rangle_{u, k_0}}$$

A much simpler protocol
(for spins)

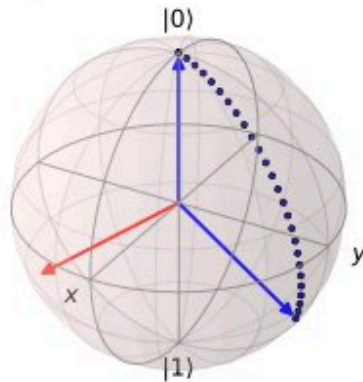
1st measurement (day 1)



2nd measurement (day 2)



Single spin random rotation

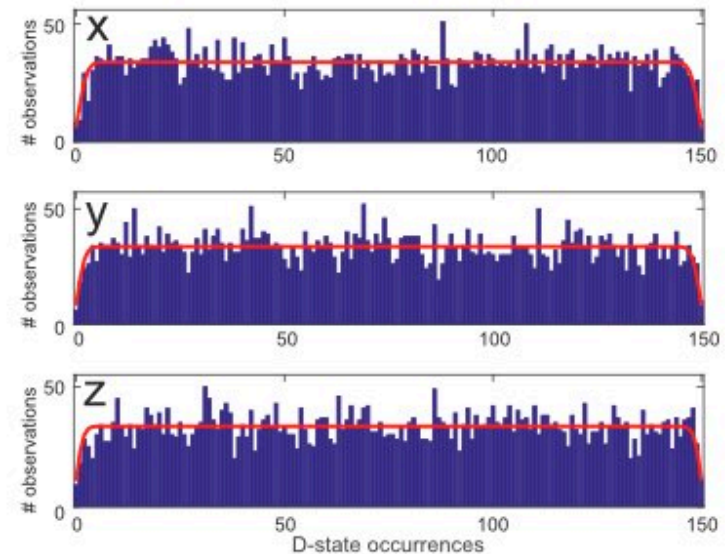


$$u_i \equiv u(\alpha_i, \beta_i, \gamma_i)$$

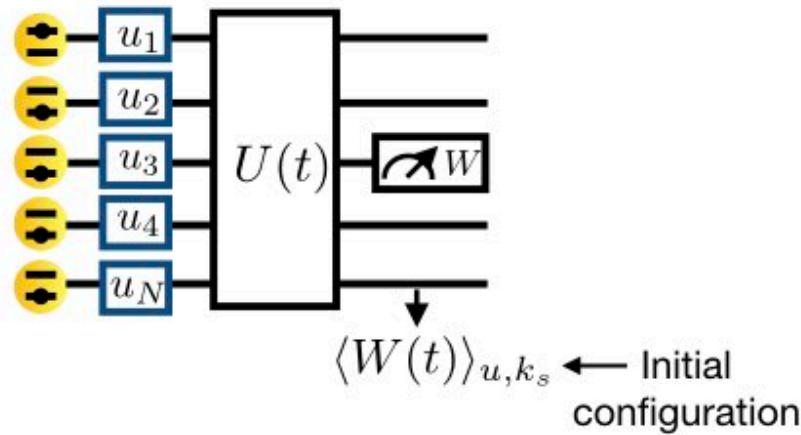
$$= Z(\alpha_i) Y(\pi/2) Z(\beta_i) Y(-\pi/2) Z(\gamma_i)$$

Available in state-of-the-art setups

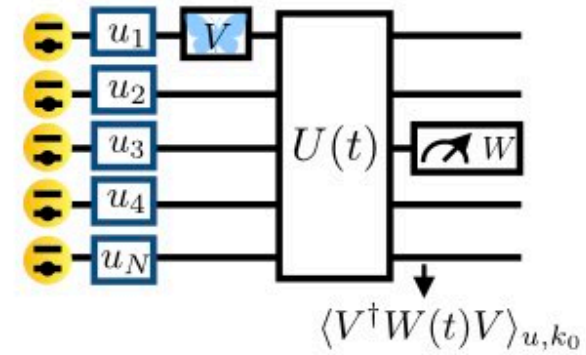
Brydges et al 2018 arXiv:1806.05747



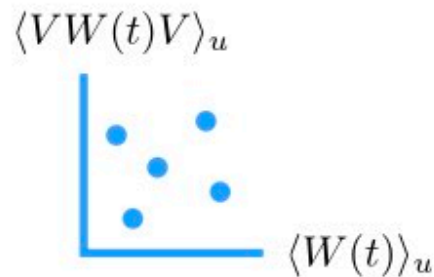
1st measurement (day 1)



2nd measurement (day 2)



Local statistical Correlations = Modified OTOCs



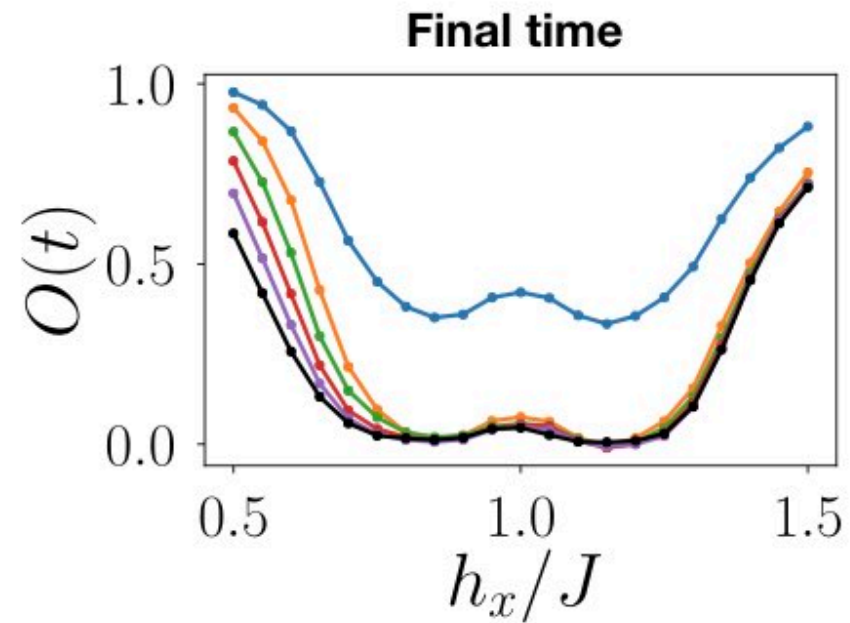
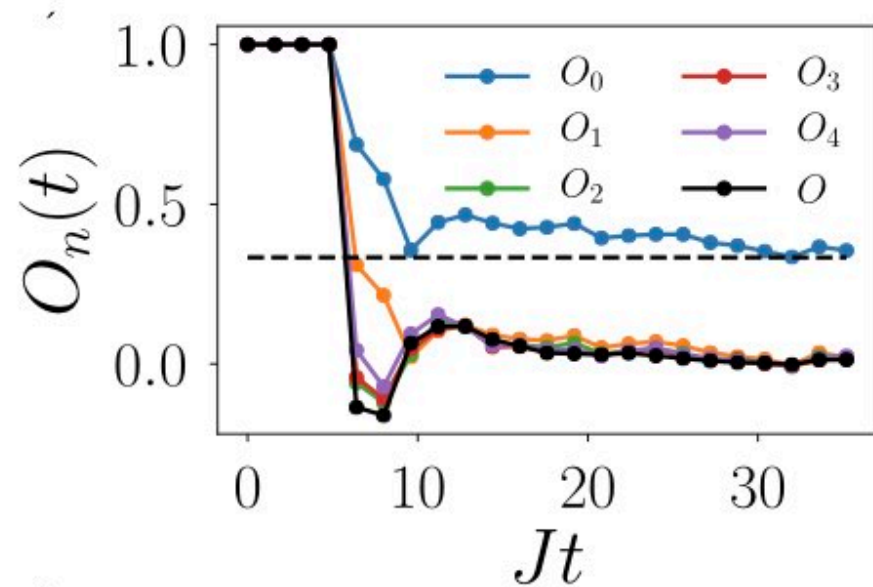
$$O_n(t) = \frac{1}{\mathcal{D}_n^{(L)}} \sum_{k_s \in E_n} c_{k_s} \overline{\langle W(t) \rangle_{u, k_s} \langle V^\dagger W(t) V \rangle_{u, k_0}},$$

\leftarrow Set of 2^n initial states, $n=0, \dots, N$

$$O_n(t) = \frac{\sum_{A, B_n \subseteq A} \text{Tr}_A (W(t)_A (VW(t)V)_A)}{\sum_{A, B_n \subseteq A} \text{Tr}_A (W(t)_A W(t)_A)}$$

- Modified OTOCs:**
- $\rightarrow O_N(t) = O(t)$
 - \rightarrow Fast converging series: $n=0, 1, 2$ is generically sufficient

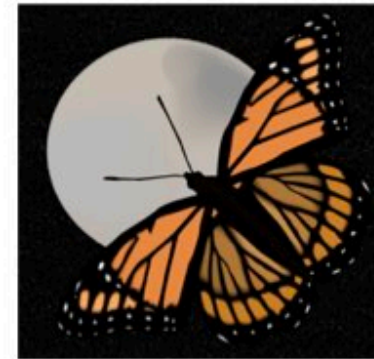
Example of Many-body Chaos: Kicked Ising with 8 sites



OTOCs as diagnosis of scrambling

[arxiv:1807.09087](https://arxiv.org/abs/1807.09087)

can be measured in many-body systems with current technology



AMO implementations (no copies)

Statistical errors are not a fundamental issue

Natural robustness against errors and imperfections (ex: depolarization)

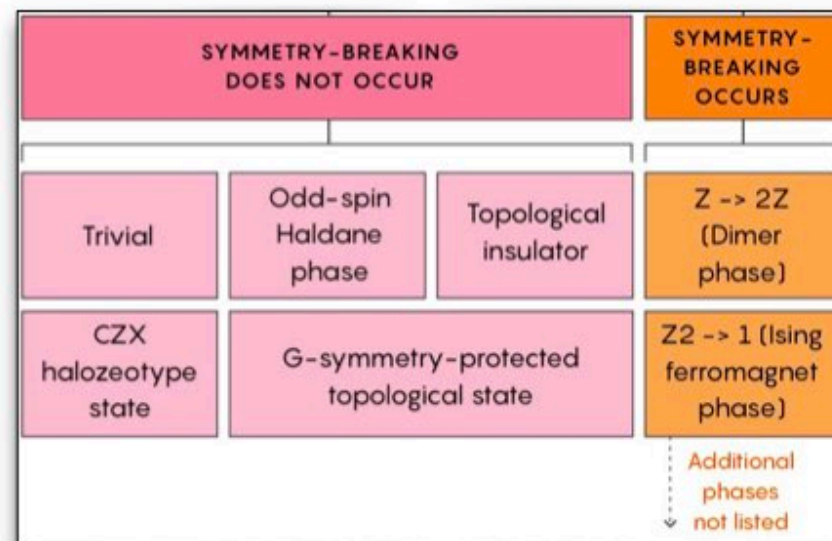
Experiment in progress (M. Joshi, T. Brydges, C. Maier, C. Roos, and R. Blatt)

Measuring scrambling with random measurements

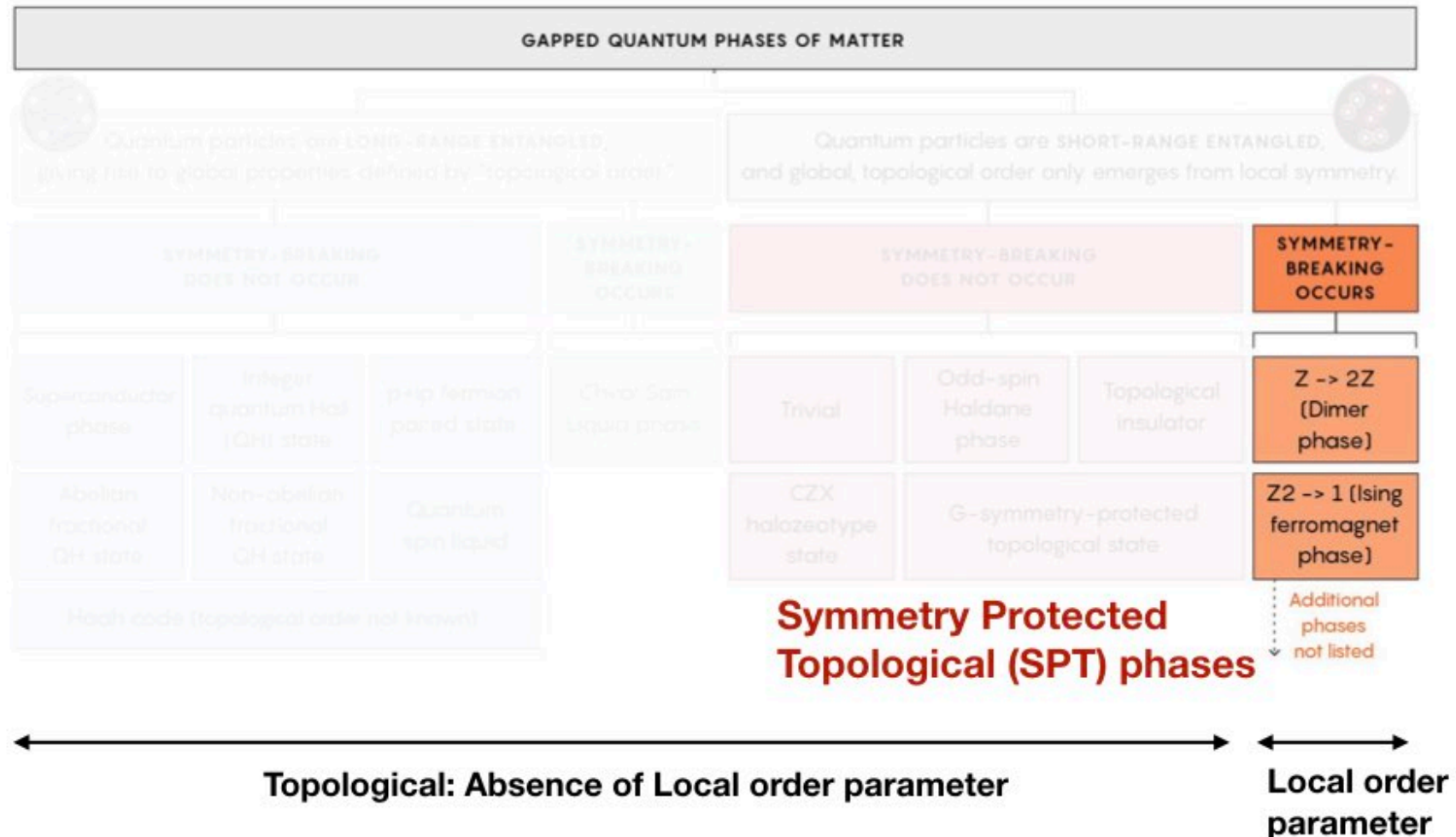


B. Vermersch, A. Elben, L. Sieberer,
N. Yao, and P. Zoller

Classification of interacting topological phases (SPT)



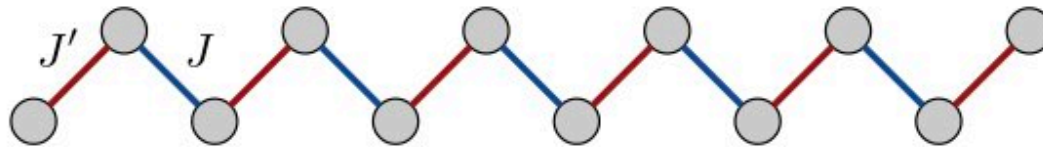
A. Elben, B. Vermersch, J. Yu, G. Zhu,
M. Hafezi and P. Zoller



Pollmann et al. PRB 2010, Schuch et al. PRB 2011, Chen et al., Science 2012, PRB 2013, ...

How to probe their classification in the lab?

1D spin-1/2 model with alternating hoppings



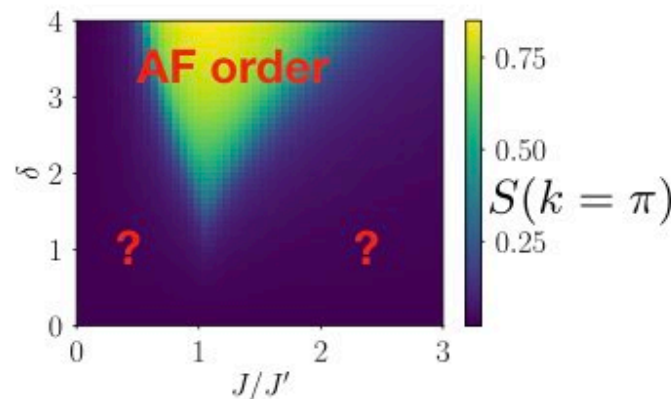
$$H = J' \sum_{i=1}^N \left(\sigma_{2i-1}^- \sigma_{2i}^+ + \text{h.c.} + \frac{\delta}{2} \sigma_{2i-1}^z \sigma_{2i}^z \right) + J \sum_{i=1}^{N-1} \left(\sigma_{2i}^- \sigma_{2i+1}^+ + \text{h.c.} + \frac{\delta}{2} \sigma_{2i}^z \sigma_{2i+1}^z \right)$$

SB phase $|J'| \approx |J|, \delta \gg 1$



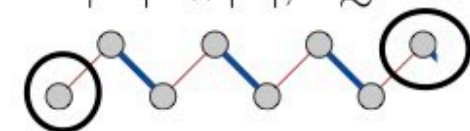
Trivial phase

$|J'| \gg |J|, \delta \lesssim 1$



Topological phase

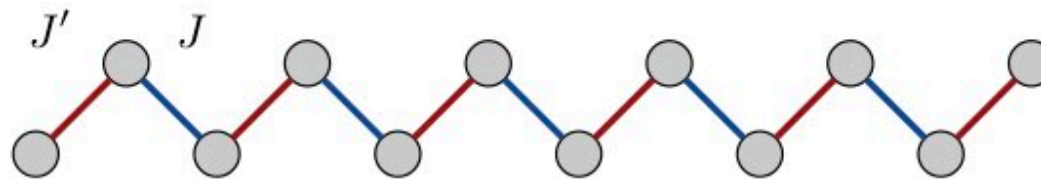
$|J'| \ll |J|, \delta \lesssim 1$



No local order parameter: **topological order**

Classification of SPT phases ↔ Classification of the action of symmetry groups

Pollmann et al. PRB 2010, Schuch et al. PRB 2011, Chen et al., Science 2012, PRB 2013, ...



Type of symmetries

- Internal symmetries (rotations)
- Bond-centered inversion
- Time-reversal

The action of symmetry groups is measured by SPT Topological invariants

Haegeman et al., PRL, 2012
Pollmann, Turner, PRB, 2012

Classification by direct identification of each symmetry representation $[e^{i\phi(g,h)}]$

Quantized

Strong Connections with topological quantum field theory, tensor-network theory

Ryu, PRL 2017

Key quantities for the classification

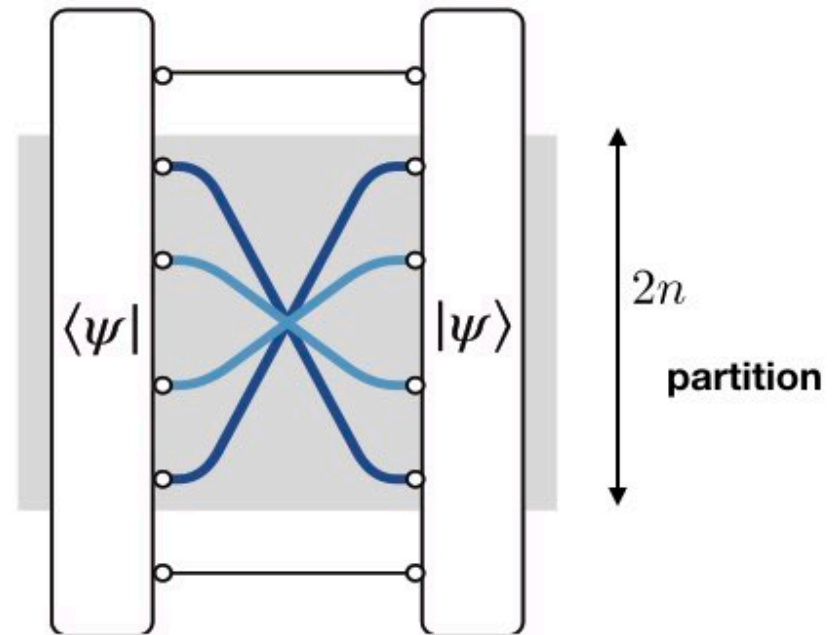
No protocols so far

Partial inversion invariant

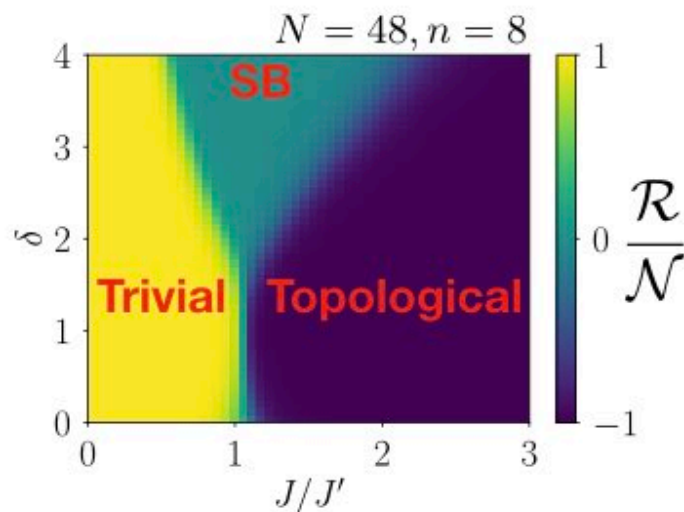
Pollmann, Turner, PRB 2012

$$\mathcal{R}(n) = \text{Tr} [\mathcal{S}_{I_1, I_2} |\Psi\rangle \langle \Psi|]$$

MPS theory
 $\xrightarrow{n \rightarrow \infty}$ $\pm \text{Tr} [\rho_{HP}^2]$ ← purity



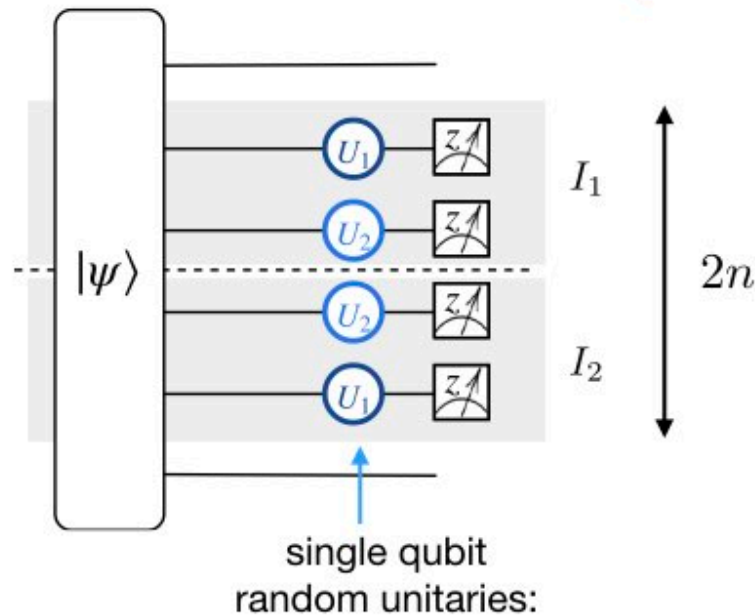
Application to the SSH model



The partial inversion invariant classifies the whole SSH model

How to measure such *non-local* correlations in an experiment?

Idea: Correlate random unitaries *in space*



$$d^n \sum_{\mathbf{s}_{I_1}, \mathbf{s}'_{I_2}} (-d)^{-D[\mathbf{s}_{I_1}, \mathbf{s}'_{I_2}]} \overline{P_{U \otimes U}(\mathbf{s}_{I_1}, \mathbf{s}'_{I_2})}$$

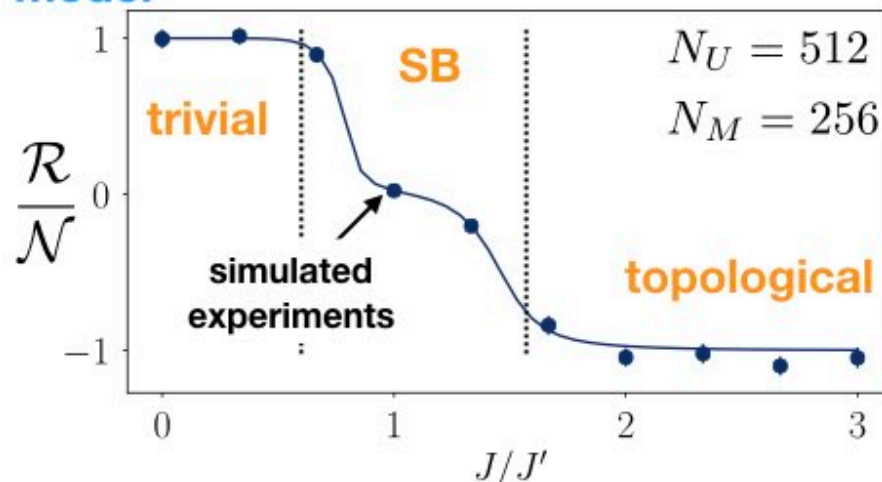
$$= \text{Tr} [\mathbf{S}_{I_1, I_2} |\Psi\rangle \langle \Psi|]$$

$$\xrightarrow{n \rightarrow \infty} \pm \text{Tr} [\rho_{HP}^2]$$

Classification via random measurements:

apply a distribution with encodes the symmetry to characterize

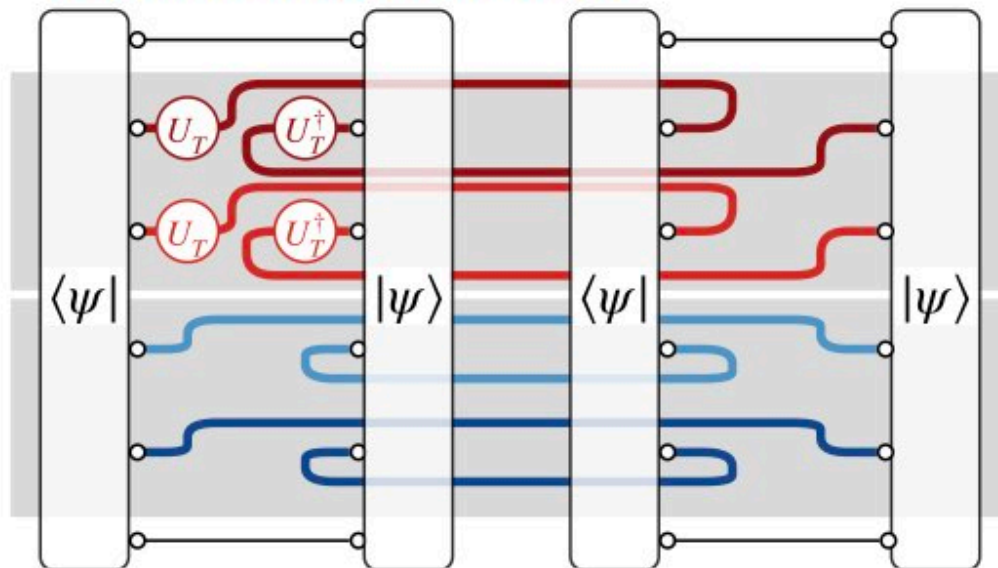
SSH - model



Error bars and bias correction with Jackknife resampling

$$N = 24, n = 6, \delta = 2.5$$

Partial transpose invariant

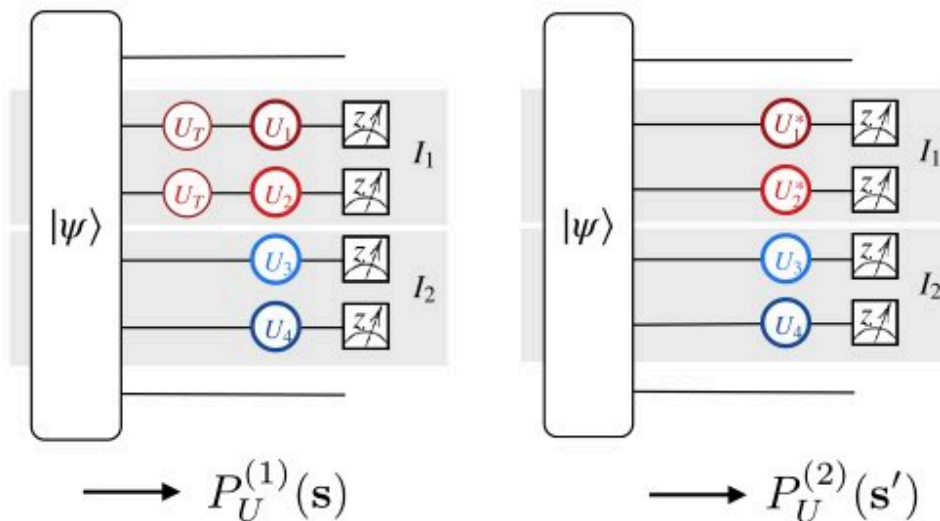


$$\mathcal{T}(n) = \text{Tr} [R_{I_1} S_{I_2} |\Psi \otimes \Psi\rangle \langle \Psi \otimes \Psi|]$$

MPS theory

$$\xrightarrow{n \rightarrow \infty} \pm \text{Tr} [\rho_{HP}^2]^3$$

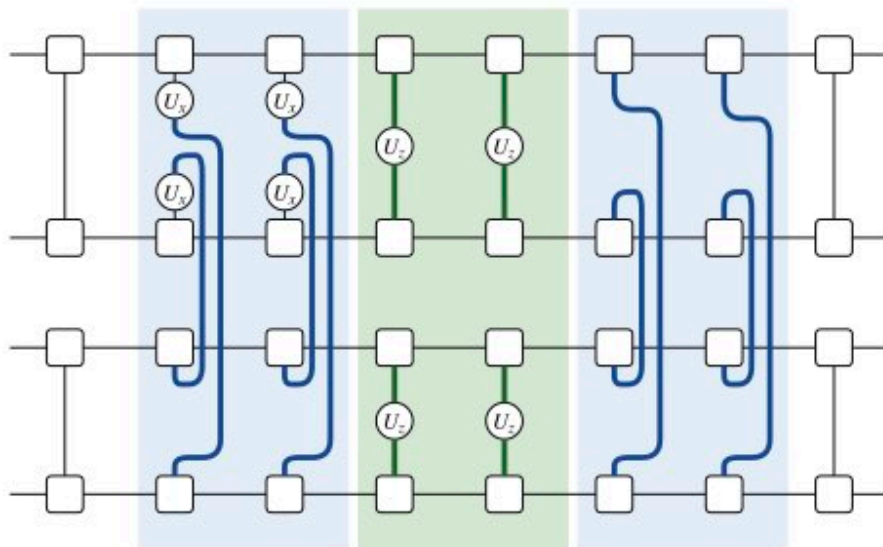
Protocol: Correlate two experiments



$$d^{2n} \sum_{\mathbf{s}, \mathbf{s}'} (-d)^{-D[\mathbf{s}, \mathbf{s}']} \overline{P_U^{(1)}(\mathbf{s})} P_U^{(2)}(\mathbf{s}') \\ = \text{Tr} [R_{I_1} S_{I_2} |\Psi \otimes \Psi\rangle \langle \Psi \otimes \Psi|]$$

Onsite unitary symmetry

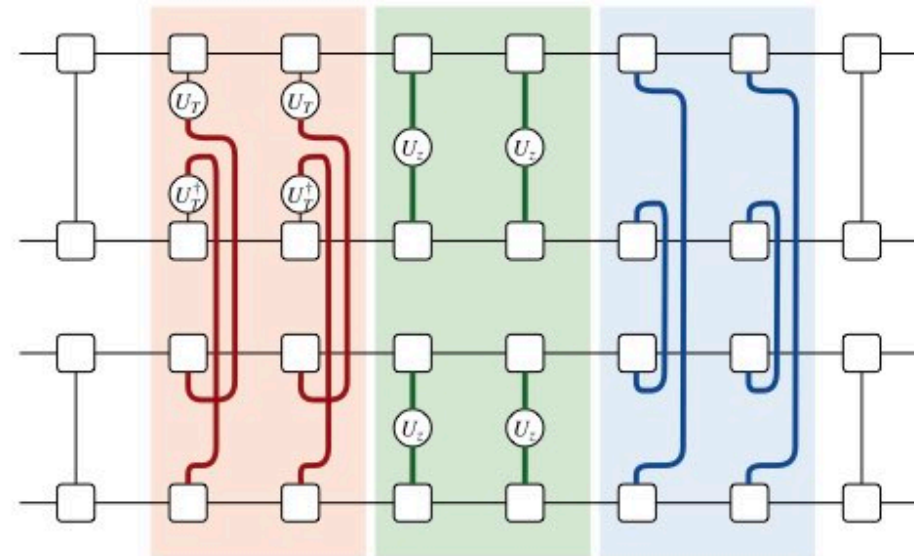
SSH: $\mathbb{Z}_2 \times \mathbb{Z}_2$ $e^{i\pi/2\sigma_x}, e^{i\pi/2\sigma_z}$



Haegemann et al., PRL 2012

Time reversal + Onsite symmetry

SSH: Time reversal + $U(1)$



Shiozaki, Ryu, JHEP 2017

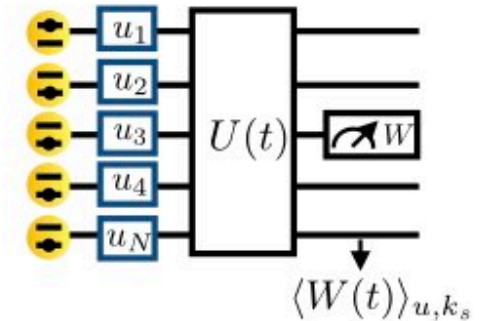
are accessible with a specific distribution of *local* random unitaires

Direct applications:

Bosonic SPT phases can be classified and tested in the lab now

Statistical correlations of randomized measurements

- a tool to probe quantum states beyond standard observables
- applicable in any state-of-the-art quantum simulation platform with high repetition rate
- **A tool to verify the quantum features of quantum simulators**



Rényi entropies

Elben, Vermersch et al. PRL, PRA 2018
 Brydges, Elben et al., arXiv:1806.05747

Out-of-time ordered correlation functions

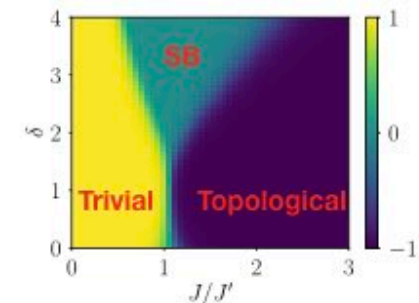
Vermersch et al., arXiv:1807.09087

Topological invariants

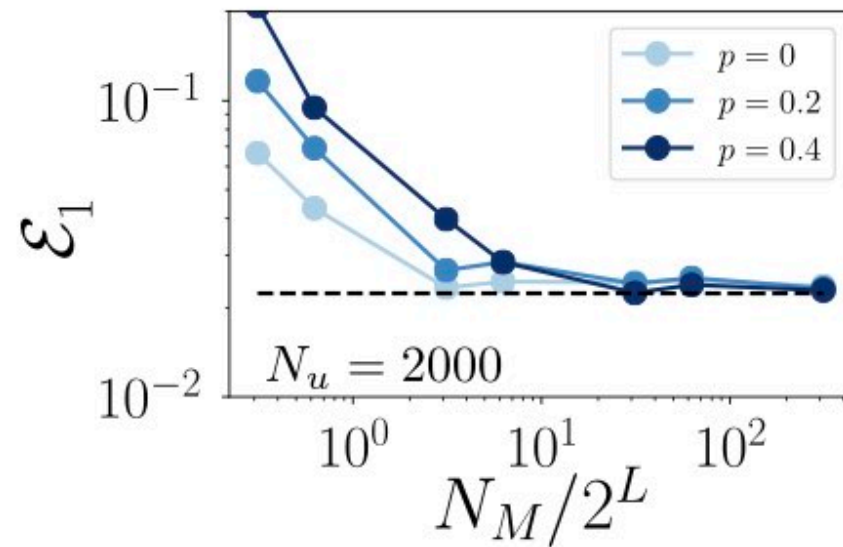
with A.Elben, J. Yu, G. Zhu,
 M. Hafezi and P. Zoller

Prospects

- Detection/Classification of true topological order
- Protocols for Hubbard models (MBL as resource?)
- Theory of random measurements
- Verification

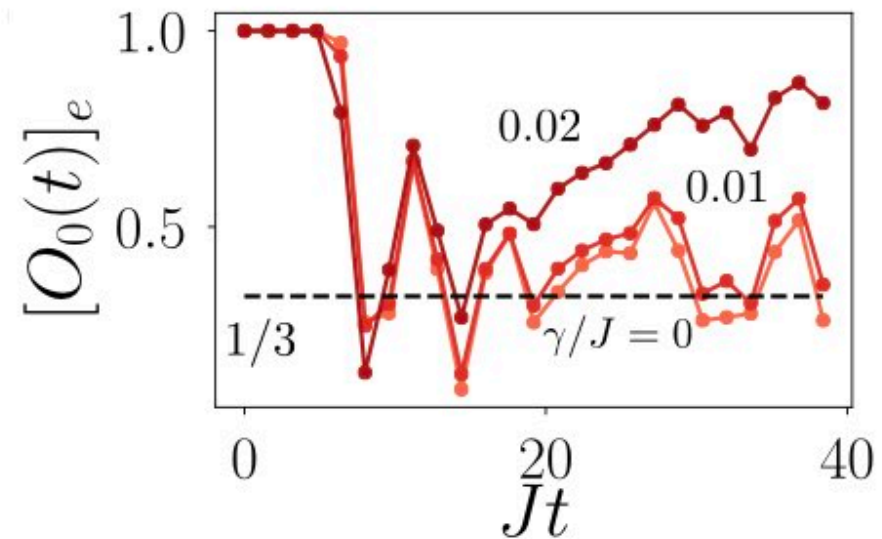


Robust against local unitary errors depolarization



Spontaneous emission

In contrast to time-reversal methods,
decoherence and scrambling
have opposite signatures

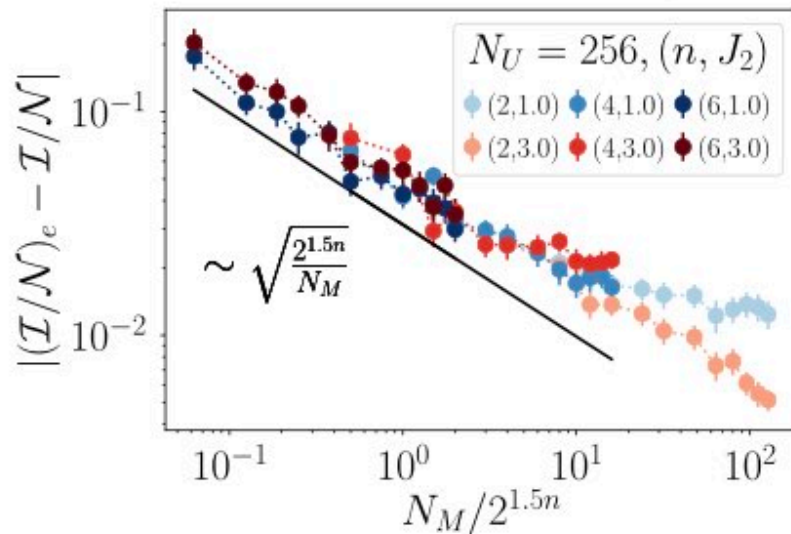


Average statistical errors

How does the required number of measurements scale with swapped sites n ?

Inversion invariant

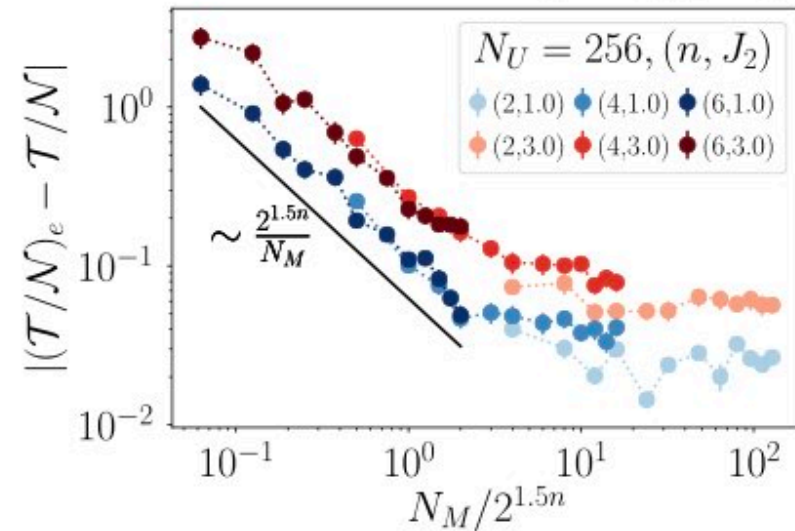
$N = 24, \delta = 0$



$$\Delta_{\mathcal{I}_n} \sim \frac{1}{\sqrt{N_U}} \left(C_I(n) + \sqrt{\frac{2^{1.5n}}{N_M}} \right)$$

Time reversal invariant

$N = 24, \delta = 0$

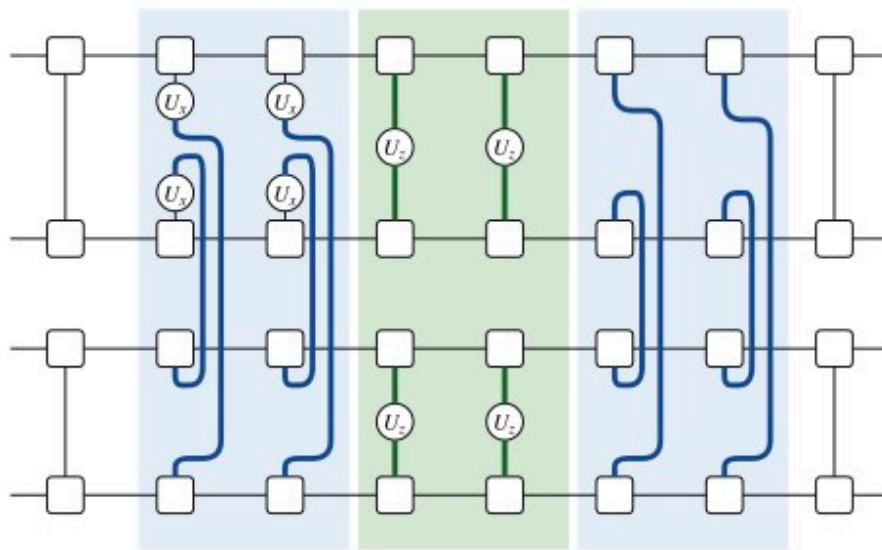


$$\Delta_{\mathcal{T}_n} \sim \frac{1}{\sqrt{N_U}} \left(C_T(n) + \frac{2^{1.5n}}{N_M} \right)$$

Exponential scaling with swapped sites - for relevant sizes within range of experimental possibilities (comparable to Renyi experiments)!

Onsite symmetry

Order parameter



$$C_n = \langle \Psi \otimes \Psi | \bigotimes_{i \in I_1} (U_i^g)^\dagger \bigotimes_{i \in I_2} (U_i^h)^\dagger S_{I_1} S_{I_3} \bigotimes_{i \in I_1} U_i^g \bigotimes_{j \in I_2} U_j^h | \Psi \otimes \Psi \rangle$$

Haegemann et al., PRL 2012

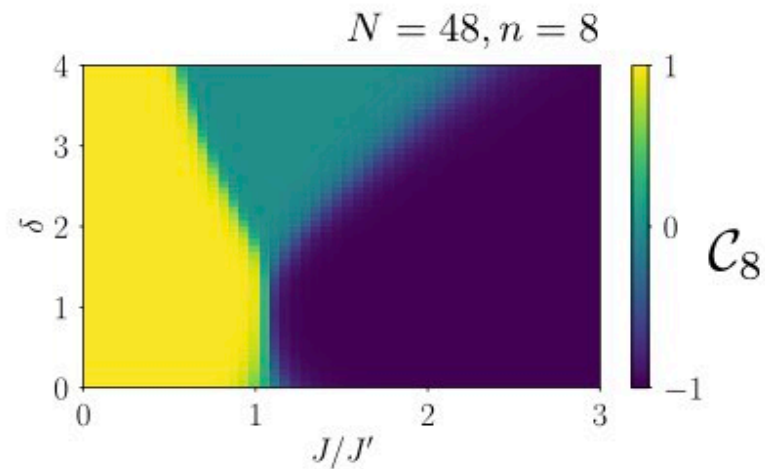
SSH model

Symmetry

$$D_2 = \mathbb{Z}_2 \times \mathbb{Z}_2$$

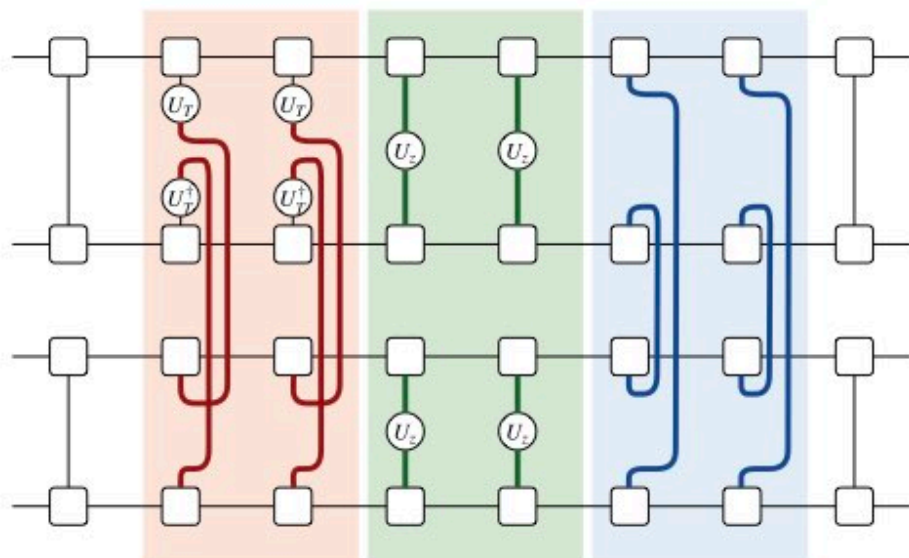
acting through

$$e^{i\pi S_x} \quad e^{i\pi S_z}$$



Klein-Bottle invariant

Time reversal + Onsite symmetry



Shiozaki, Ryu, JHEP 2017

$$\mathcal{K}_n = \langle \Psi \otimes \Psi | \bigotimes_{i \in I_1} (U_i^T)^\dagger \bigotimes_{i \in I_2} (U_i^g)^\dagger$$

$$R_{I_1} S_{I_3} \bigotimes_{i \in I_1} U_i^T \bigotimes_{j \in I_2} U_j^g | \Psi \otimes \Psi \rangle$$

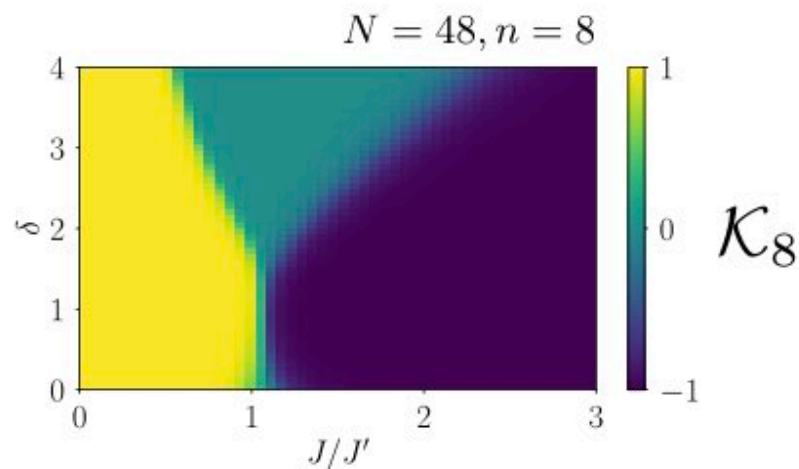
SSH model

Symmetry

$$D_2 = \mathbb{Z}_2 \times \mathbb{Z}_2 \quad + \text{Time reversal}$$

acting through

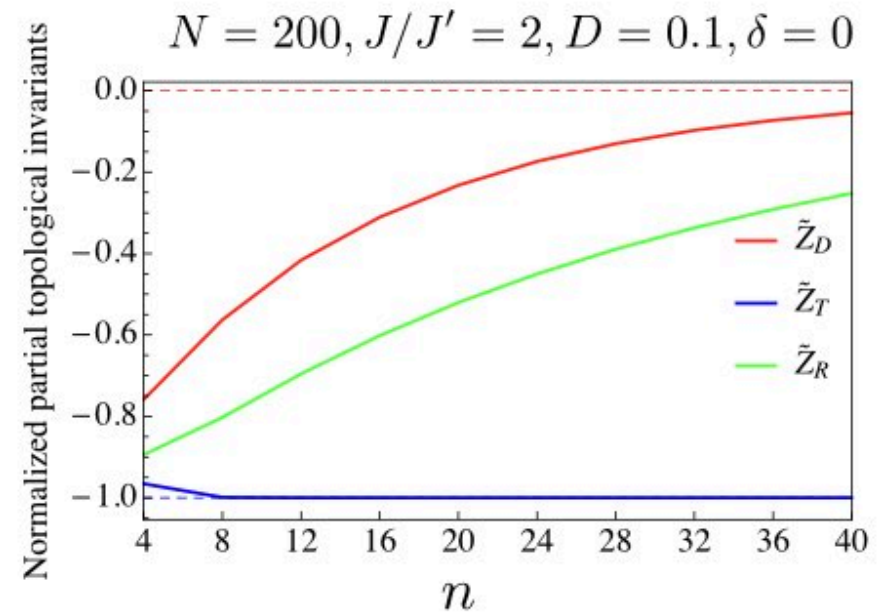
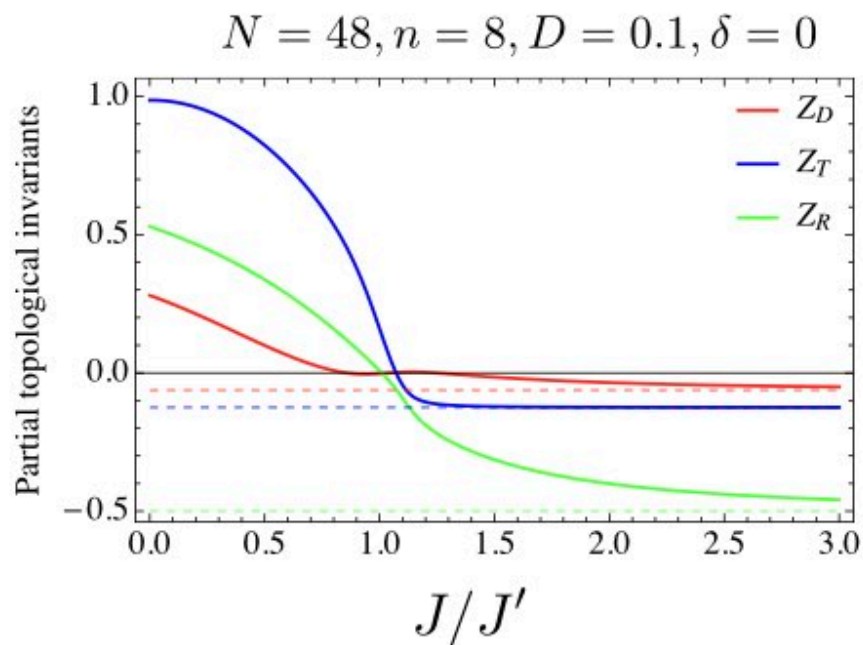
$$e^{i\pi/2\sigma_y} \mathcal{K}, e^{i\pi/2\sigma_x}, e^{i\pi/2\sigma_z}$$



Explicit breaking of symmetries

$$H = H_{\text{SSH}} + D \sum_j (\sigma_j^x \sigma_{j+1}^z - \sigma_j^z \sigma_{j+1}^x)$$

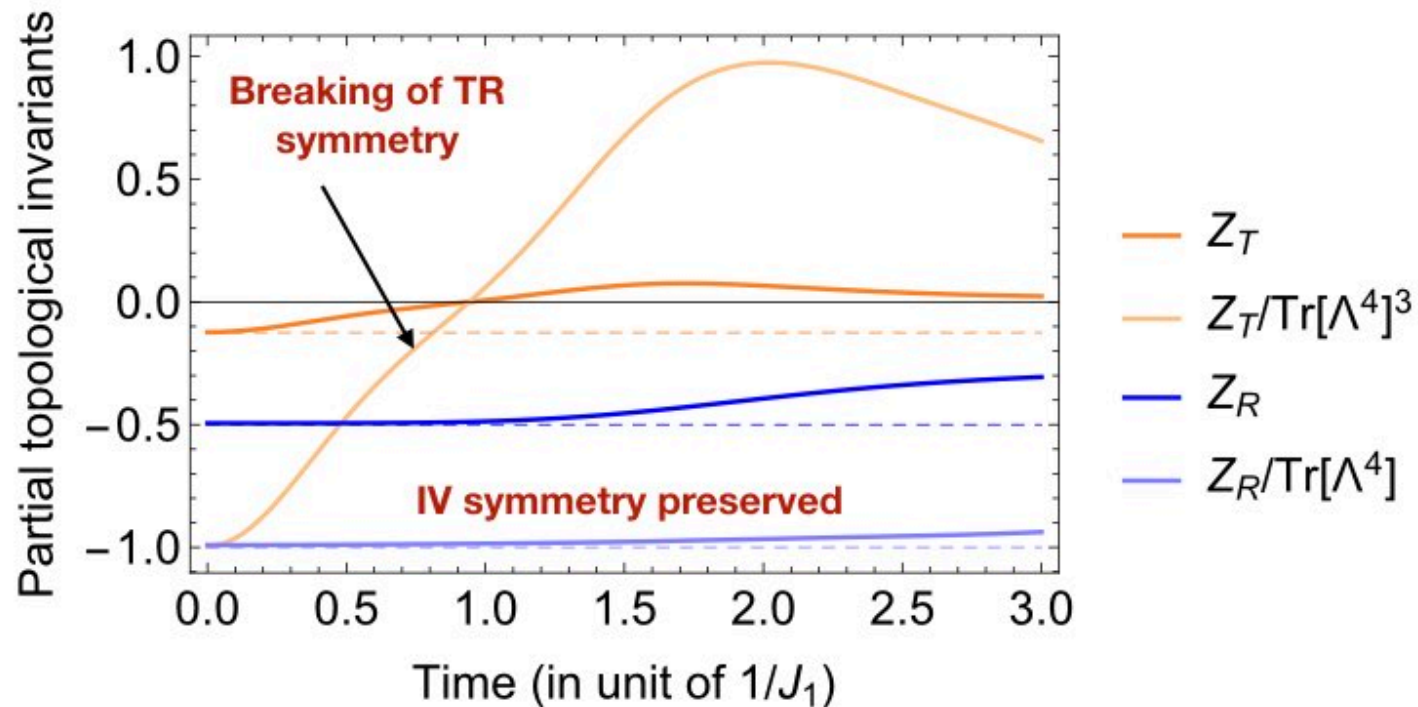
← Inversion and D2
symmetry broken



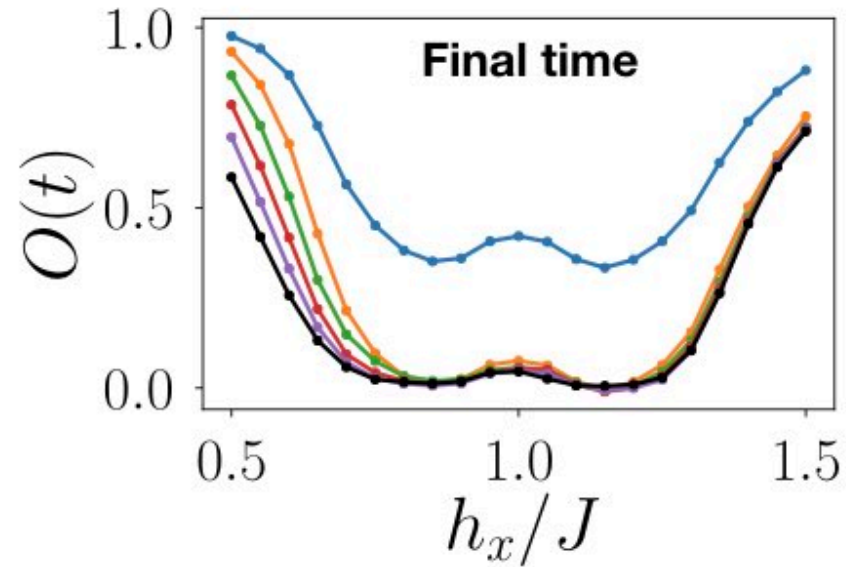
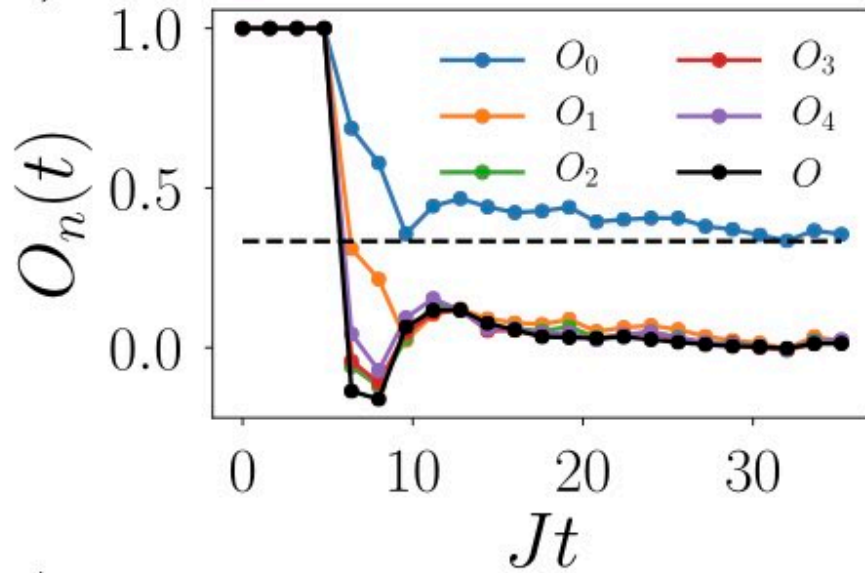
Quench dynamics - Breaking time reversal symmetry

$$H = J' \sum_{i=1}^N (\sigma_{2i-1}^- \sigma_{2i}^+ + \text{h.c.} + \delta \sigma_{2i-1}^z \sigma_{2i}^z) + J \sum_{i=1}^{N-1} (\sigma_{2i}^- \sigma_{2i+1}^+ + \text{h.c.} + \delta \sigma_{2i}^z \sigma_{2i+1}^z)$$

Quench from topological to trivial phase: $(J', J) = (1, 2) \rightarrow (1, 1/2)$



Example of Many-body Chaos: Kicked Ising with 8 sites



Statistical errors

$$L(t) \approx v_B t \text{ "Scrambling length"}$$

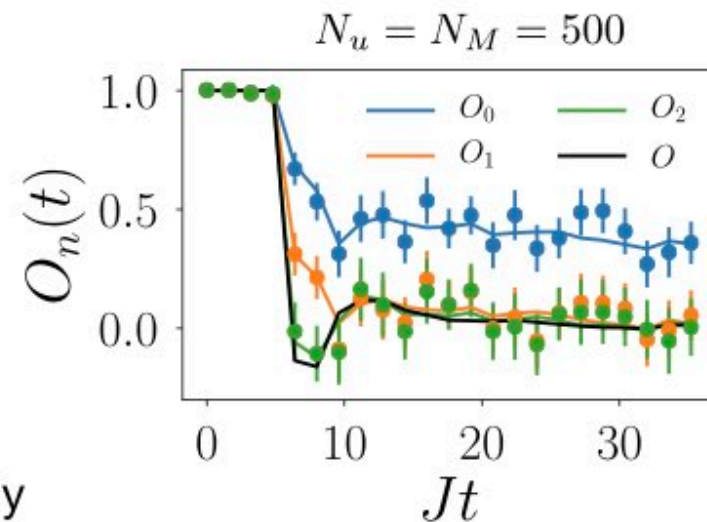
$$N_M = 2^{L(t)} \rightarrow \text{error } 1/\sqrt{N_u}$$

projective measurements

unitaries

Independent of System Size

All N operators W are measured simultaneously

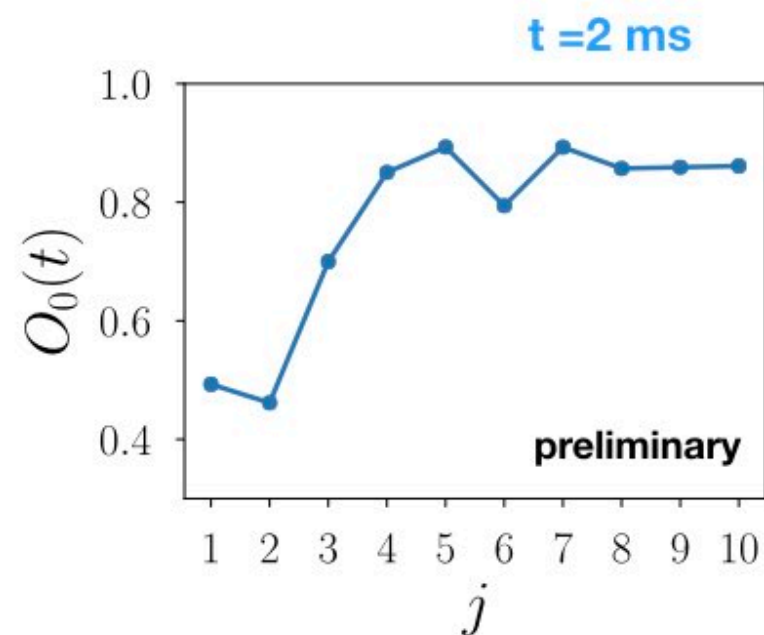
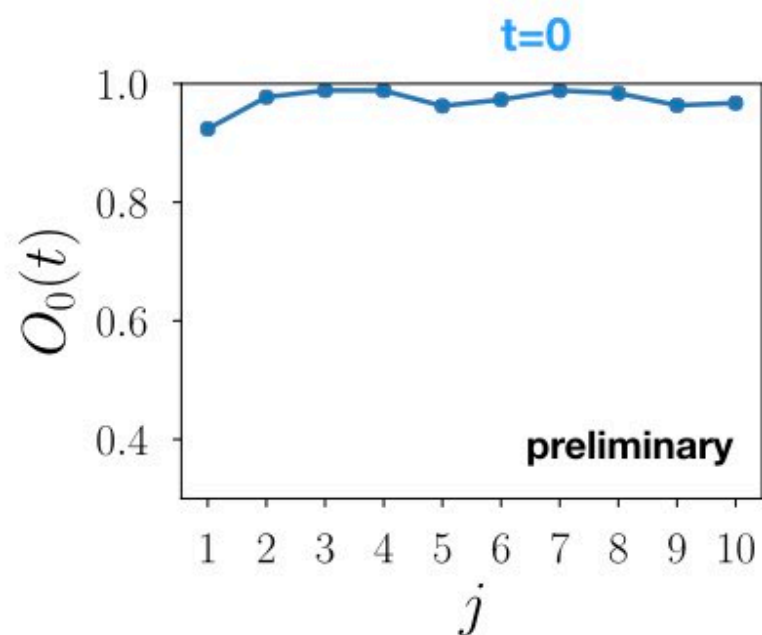


First experimental pictures

Collaboration with **M. Joshi, T. Brydges, C. Maier, C. Roos, and R. Blatt**

10 ions, evolution with long-range XY model

OTOCs in the x-basis



Haegeman et al., PRL, 2012
Pollmann et al. PRB 2012

IMPS: $\bigotimes_i U_g^i |\psi_{GS}\rangle = \text{---} \begin{array}{c} \boxed{A} \\ \circledast U_g \end{array} \text{---} \begin{array}{c} \boxed{A} \\ \circledast U_g \end{array} \text{---} \begin{array}{c} \boxed{A} \\ \circledast U_g \end{array} \text{---} \begin{array}{c} \boxed{A} \\ \circledast U_g \end{array} \text{---} \begin{array}{c} \boxed{A} \\ \circledast U_g \end{array} \text{---} = |\psi_{GS}\rangle$

Transformation under symmetry action

$$\text{---} \begin{array}{c} \boxed{\tilde{A}} \\ \circledast U_g \end{array} \text{---} = \text{---} \begin{array}{c} \diamond V_g^\dagger \\ \circledast U_g \end{array} \begin{array}{c} \boxed{A} \\ \circledast U_g \end{array} \begin{array}{c} \diamond V_g \\ \circledast U_g \end{array} \text{---}$$

with

Projective representations of G

$$V_g \tilde{V}_h = e^{i\phi(g,h)} V_{gh}$$

On-site unitary: $\tilde{A} = A$
Inversion: $\tilde{A} = A^T$
Time reversal: $\tilde{A} = A^*$

Classification of SPT phases

$$\left[e^{i\phi(g,h)} \right] \in H^2(G, U(1)_\phi)$$

← 2nd cohomology group of G

Chen et al., Science 2012

How to access and measure $\left[e^{i\phi(g,h)} \right]$ for a given G?

Purity from random measurements

Observation 1

$$\text{Tr} [\rho^2] = \text{Tr} [\mathbb{S} \rho \otimes \rho] \quad \text{with Swap operator } \mathbb{S}$$

Observation 2

$$\sum_{\mathbf{s}, \mathbf{s}'} A_{\mathbf{s}, \mathbf{s}'} \overline{P_U(\mathbf{s}) P_U(\mathbf{s}')} = \text{Tr} \left[\underbrace{U^\dagger \otimes U^\dagger A U \otimes U}_{\Phi_N^2(A)} \rho \otimes \rho \right]$$

$$P_U(\mathbf{s}) = \text{Tr} [U \rho U^\dagger |\mathbf{s}\rangle \langle \mathbf{s}|]$$

$$U = \bigotimes u_i \quad u_i \in \text{CUE}(d)$$

Generalized unitary twirling channel on 2-copy space

$$\xrightarrow{\text{Schur-Weyl duality}} \Phi_N^2(A) = \sum_{\pi, \sigma \in S_2^{\otimes N}} C_{\pi, \sigma} \text{Tr} [W_\pi A] W_\sigma \stackrel{!}{=} \mathbb{S}$$

↑ Weingarten Matrix
↑ Permutation operator

$$\longrightarrow A = 2^N \sum_{\mathbf{s}, \mathbf{s}'} (-2)^{-D[\mathbf{s}, \mathbf{s}]} |\mathbf{s}\rangle \langle \mathbf{s}| \otimes |\mathbf{s}'\rangle \langle \mathbf{s}'|$$