

# A New Solid-State Platform for Quantum Simulation and Optimization

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[arXiv:1809.03794](https://arxiv.org/abs/1809.03794)

# Motivation: Quantum Optimization (QAOA)

Motivation: Clustering (unsupervised learning).

$$w_{i,j} = d(\mathbf{x}_i, \mathbf{x}_j)$$

Maximization goal:

$$\sum_{i \in S, j \in \bar{S}} w_{i,j} \rightarrow \max$$

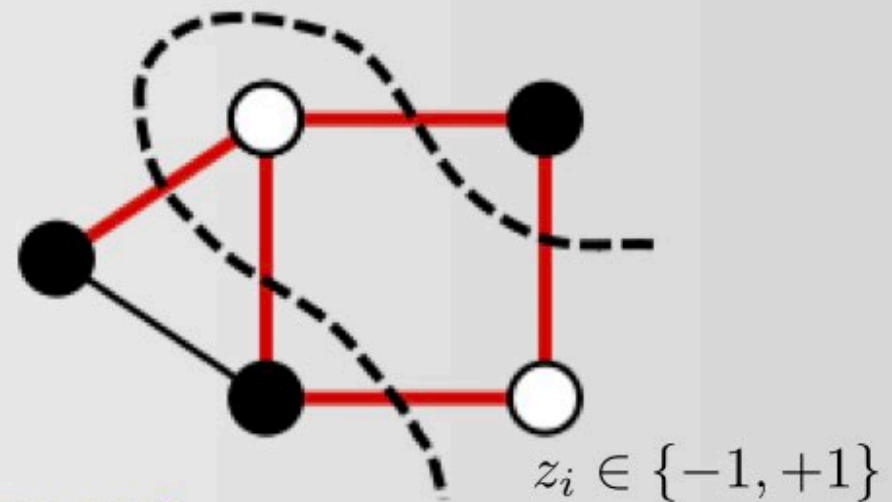
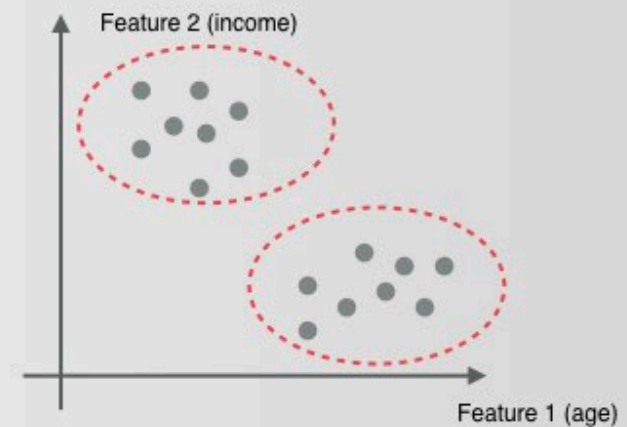
Put the problem on a graph:

$$H_C = \sum_{i,j} \frac{w_{i,j}}{2} (1 - z_i z_j)$$

[MAXCUT: NP-complete problem]

Minimize energy of frustrated anti-ferromagnet:

$$H_{AF} = \sum_{i,j} w_{i,j} z_i z_j$$



# Quantum Optimization (QAOA)

## Quantum Adiabatic Algorithm:

- (Classical) problem Hamiltonian:  $H_C = \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z$
- **Goal:** Find/approximate ground state (energy)

$$H(s) = s \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z + (1-s) \sum_i \sigma_i^x \quad s = t/T$$

$$s = 0, H = \sum_i \sigma_i^x$$

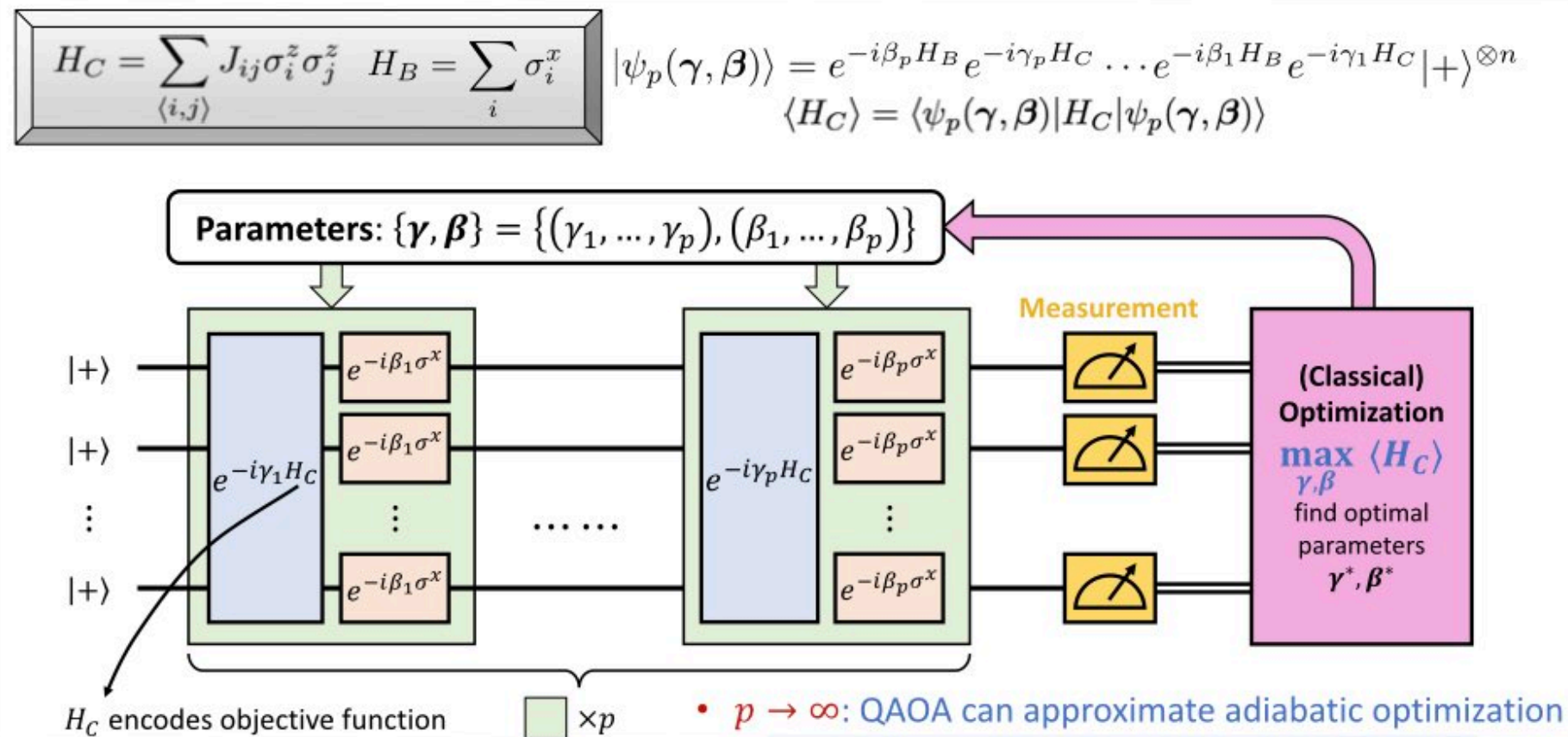
$$s = 1, H = \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z$$



- Gap can be exponentially small (with the system size)

# Quantum Optimization (QAOA)

## Quantum Approximate Optimization Algorithm (QAOA):



arXiv:1411.4028

- QAOA: Heuristic *variational* ansatz (Edi Farhi, MIT).
- Apply iteratively pair of unitaries before measuring classical string.
- Probability of success approaches unity for  $M \rightarrow \infty$ .



# Requirements for a solid-state QAOA architecture

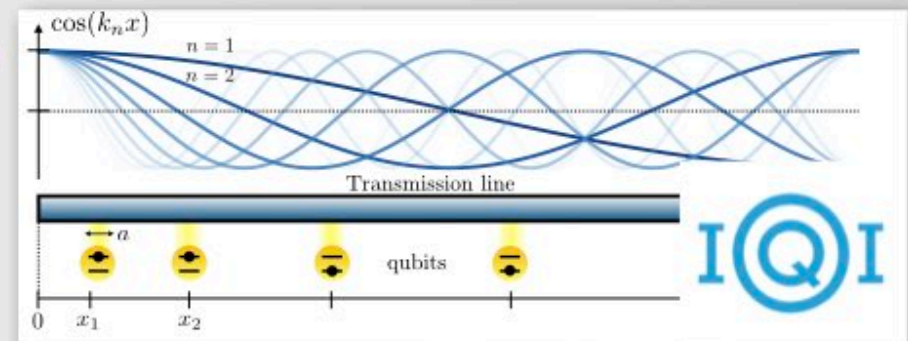
Single qubit rotations

$$H_C = \sum_{\langle i,j \rangle} J_{ij} \sigma_i^z \sigma_j^z \quad H_B = \sum_i \sigma_i^x$$

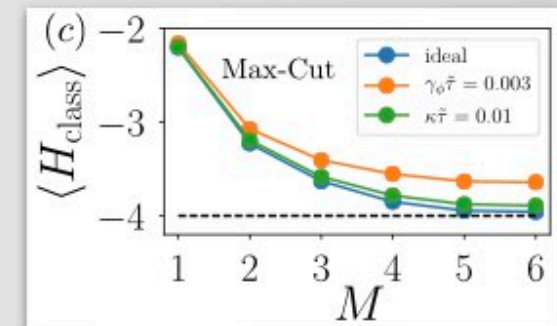
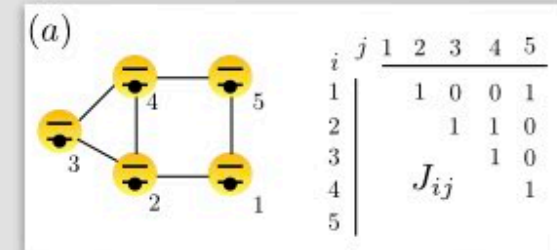
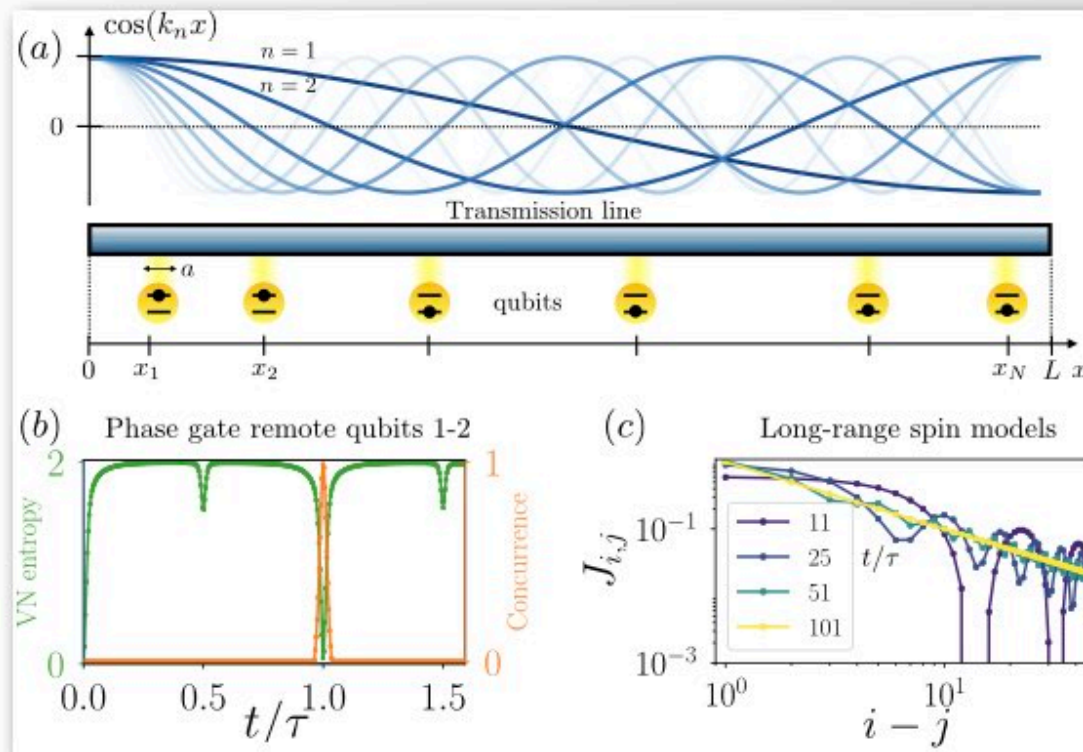
Arbitrary All to All Ising interactions

- **Programmable**
- Scalable
- Non-Perturbative
- Robust

- Quantum Simulation and Optimization in Hot Quantum Networks:
  - Flexible and robust platform for QAOA.
  - arXiv: 1809.03794.



# Take-Home Message



## Long-range entanglement and quantum optimization in hot quantum networks:

- Long (*multi-mode*) transmission line as quantum bus.
- No ground-state cooling required, data bus can be **hot**.
- Approach not based on perturbative argument for qubit-photon coupling.
- Qubits do *not* have to be identical.
- Recipe to generate desired spin-spin interactions.
- Robust and flexible implementation for QAOA.

## **Model and exact Solution**

### **Illustrations**

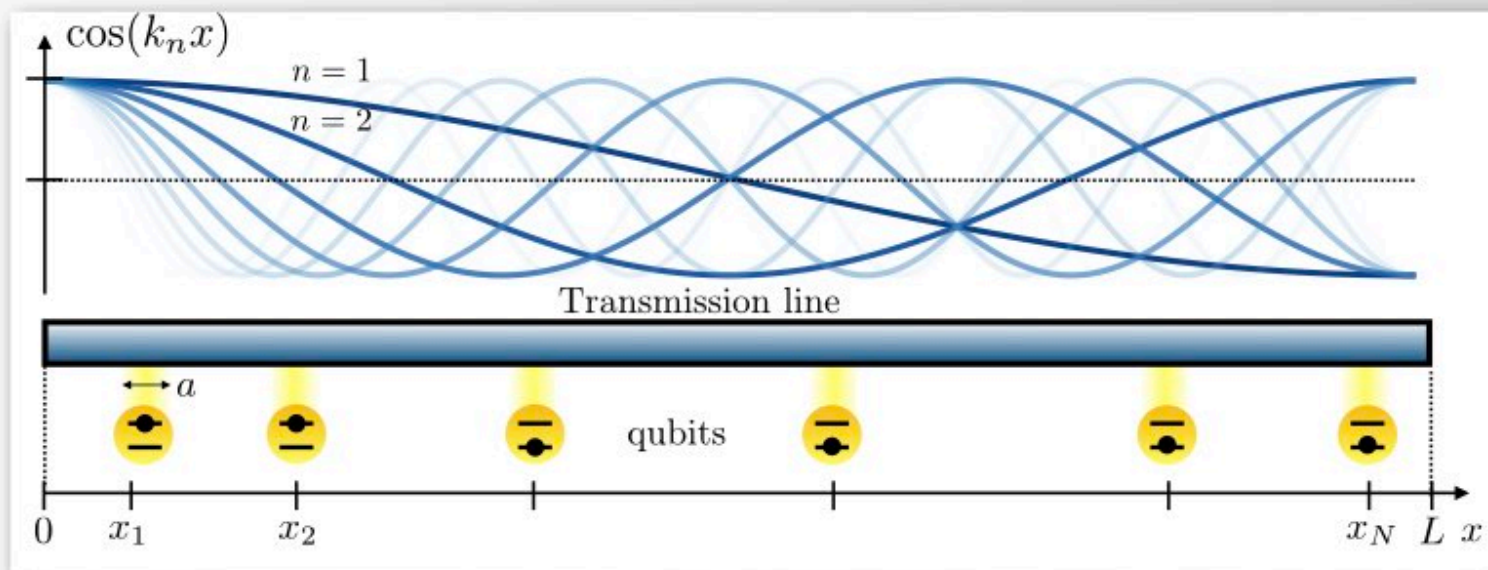
Two-qubit gate

Quantum Simulation

Quantum optimization

## **Implementations with Superconducting qubits**



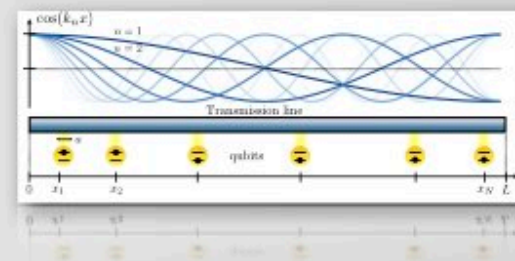


$$H = \sum_i \frac{\omega_i}{2} \sigma_i^z + \sum_{n=1}^{\infty} \omega_n a_n^\dagger a_n + \sum_{i,n} g_{i,n} \sigma_i^z (a_n + a_n^\dagger)$$

Spin-resonator system with *longitudinal* coupling:

- Linear resonator spectrum:  $\omega_n = k_n c = n\omega_1$ , where  $\omega_1 = \pi c/L$

$$H = \sum_i \frac{\omega_i}{2} \sigma_i^z + \sum_{n=1}^{\infty} \omega_n a_n^\dagger a_n + \sum_{i,n} g_{i,n} \sigma_i^z (a_n + a_n^\dagger)$$



## Multi-mode (polaron) transformation

$$H = U_{\text{pol}} \tilde{H} U_{\text{pol}}^\dagger$$

$$U_{\text{pol}}^\dagger = \exp\left[\sum_{n,i} \frac{g_{i,n}}{\omega_n} \sigma_i^z (a_n^\dagger - a_n)\right],$$

$$\tilde{H} = \sum_i \frac{\omega_i}{2} \sigma_i^z + \sum_n \omega_n a_n^\dagger a_n + \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z,$$

[exact result]

$$J_{ij} = -2 \sum_n \frac{g_{i,n} g_{j,n}}{\omega_n}.$$

[effective spin-spin interactions]

# Analytical Solution of Lab-Frame Dynamics

$$H = U_{\text{pol}} \tilde{H} U_{\text{pol}}^\dagger$$

$$\tilde{H} = \sum_i \frac{\omega_i}{2} \sigma_i^z + \sum_n \omega_n a_n^\dagger a_n + \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z,$$

$$e^{-iHt} = U_{\text{pol}} e^{-i\tilde{H}t} U_{\text{pol}}^\dagger$$

$$\omega_n t_p = 2\pi p_n \quad \text{Stroboscopic, not perturbative.}$$

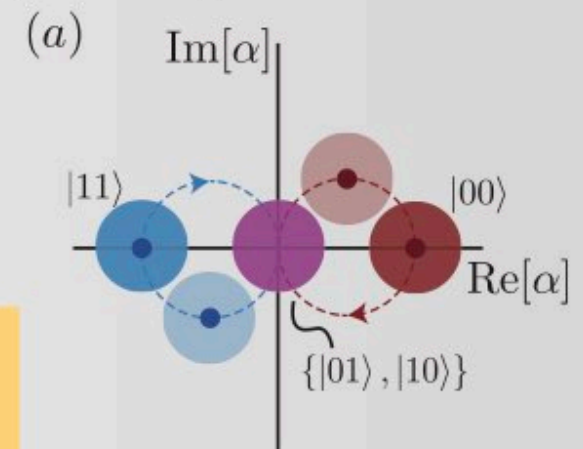
$$\exp[-it_p \sum_n \omega_n a_n^\dagger a_n] = \exp[-2\pi i \sum_n p_n a_n^\dagger a_n] = \mathbb{1}$$

$$t_p = p_1 \tau \quad (\text{with } p_n = np_1)$$

$$\tau \equiv 2L/c$$

$$U_{\text{lab}}(t_p) = e^{-it_p \sum_i (\omega_i/2) \sigma_i^z} e^{-it_p \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z}.$$

[unitary generates qubit-qubit interactions, independent of state of resonator modes]

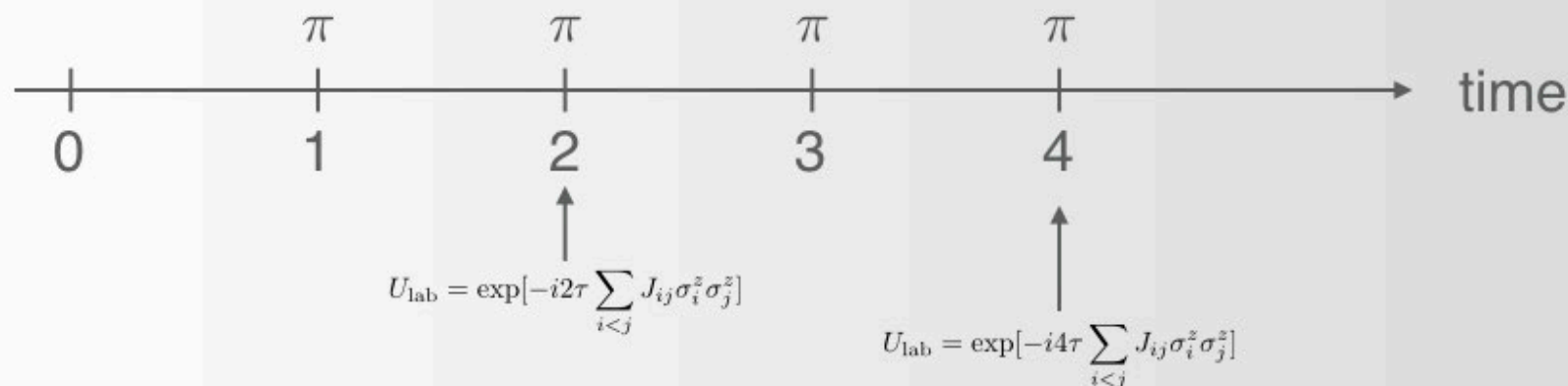


# Analytical Solution of Lab-Frame Dynamics

$$U_{\text{lab}}(t_p) = e^{-it_p \sum_i (\omega_i/2) \sigma_i^z} e^{-it_p \sum_{i<j} J_{ij} \sigma_i^z \sigma_j^z}.$$

Pure spin Hamiltonian: Independent of resonator modes.

- Approach not based on perturbative arguments.
- For certain times, evolution in polaron and lab frame fully coincide.
- Scheme insensitive to the state of the resonator, allowing for *thermally* robust gate, without need for ground-state cooling.
- No (resonance) conditions on qubit frequencies  $\omega_i$ .



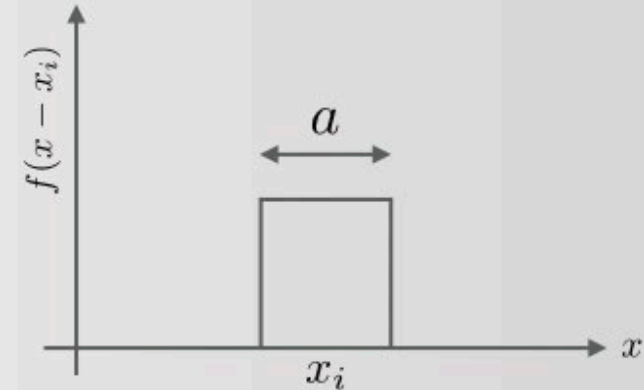
# Effective Spin-Spin Interactions: Frequency Cutoff

$$J_{ij} = -2 \sum_n \frac{g_{i,n} g_{j,n}}{\omega_n}.$$

- By definition, effective spin-spin interactions involves all modes  $n=1,2,\dots$
- Unphysical divergencies? Example:  $g_{i,n} = g_i \sqrt{n} \cos(k_n x_i)$
- Microscopic length-scale introduces frequency cutoff:  $k_{\max} \sim \pi/a$ .

$$g_{i,n} = g_i \sqrt{n} \int_0^L \cos(k_n x) f(x - x_i) dx$$

$$J_{ij} = g_i g_j / \omega_1$$



$$g_{i,n} = g_i \sqrt{n} (\sin[k_n(x_i + a)] - \sin[k_n x_i]) / (k_n a)$$



## **Model and exact Solution**

### **Illustrations**

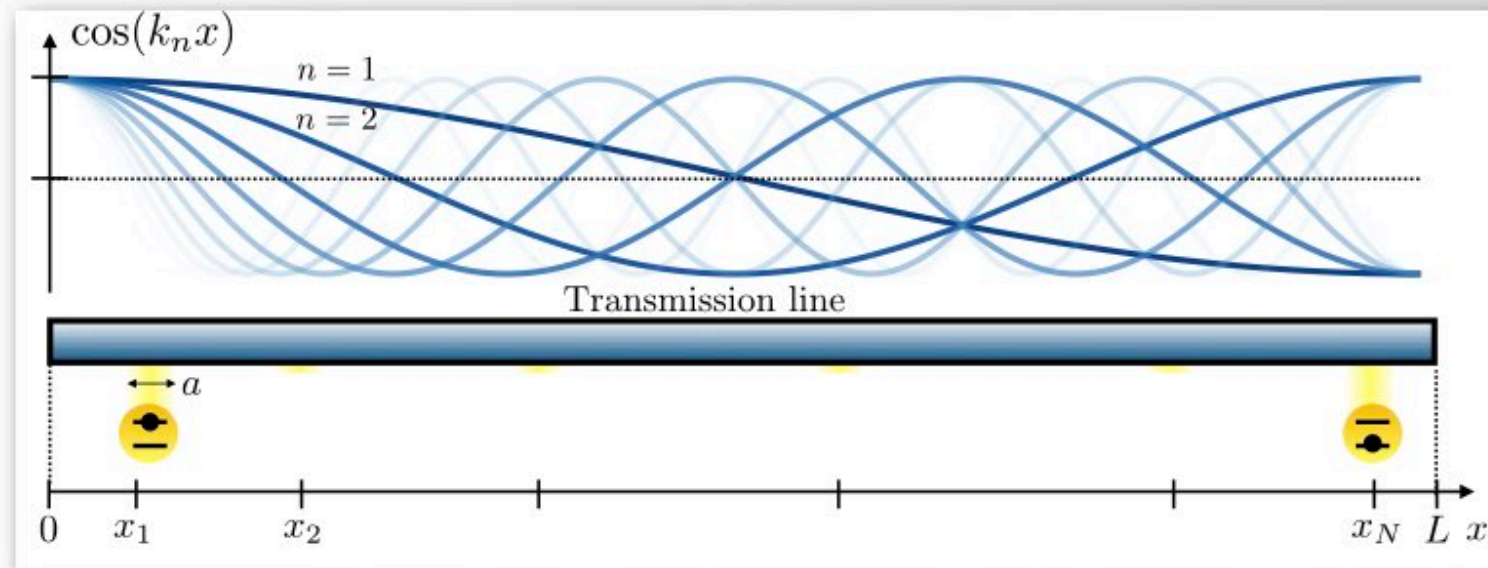
Two-qubit gate

Quantum Simulation

Quantum optimization

## **Implementations with Superconducting qubits**

## Hot qubit Gate between two remote qubits



$$\rho_0 = |\Psi_0\rangle \langle \Psi_0| \otimes_n \rho_n$$

$$|\Psi_0\rangle = \otimes_i (|0\rangle + i|1\rangle) / \sqrt{2}$$

...

$\rho_n$  : Arbitrary

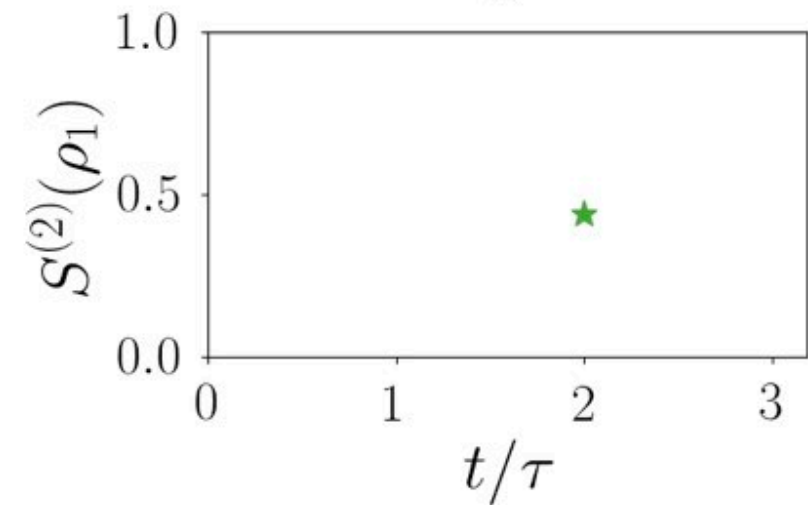
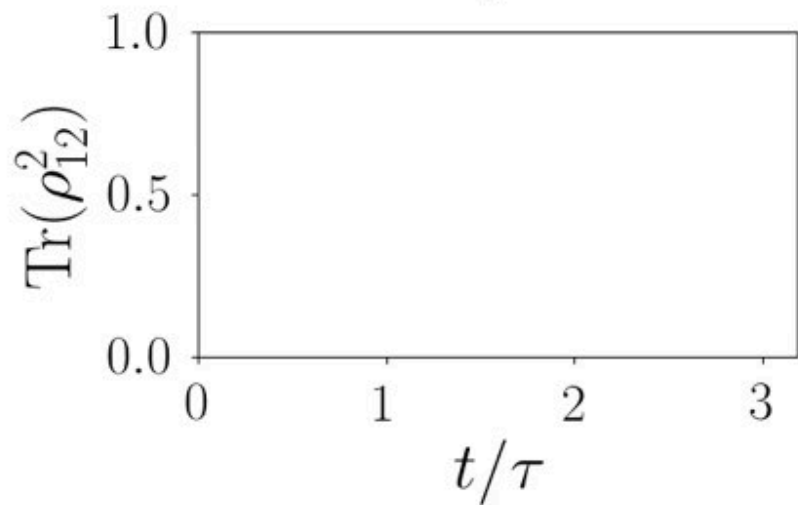
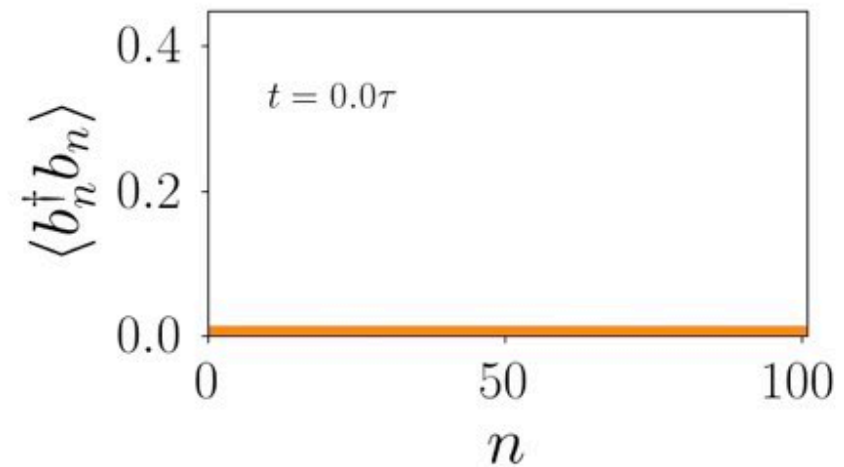
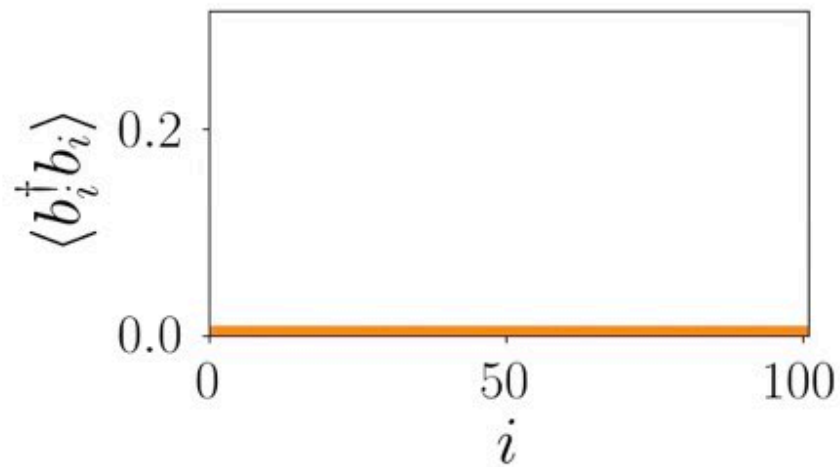
- Max. entangled target state:

$$|\Psi_{\text{target}}\rangle = \exp\left(-i\frac{\pi}{4}\sigma_1^z \sigma_2^z\right) |\Psi_0\rangle$$

# Example I

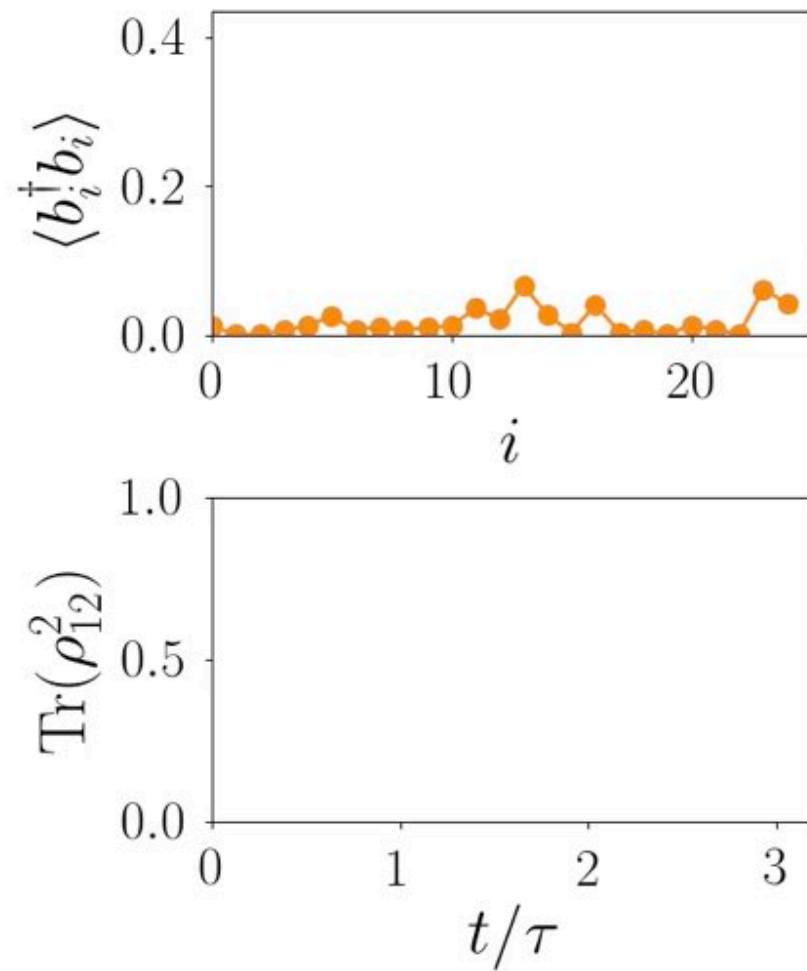
Matrix-Product State (101 Bosonic Modes)

Zero-Temperature

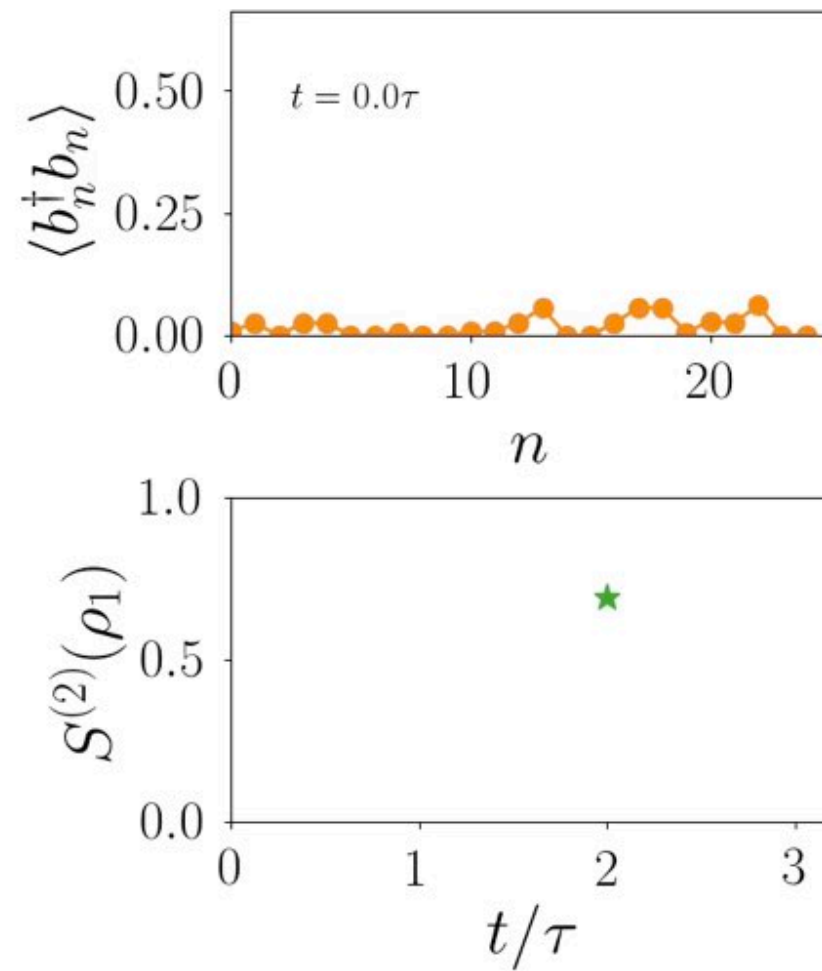


# Example I

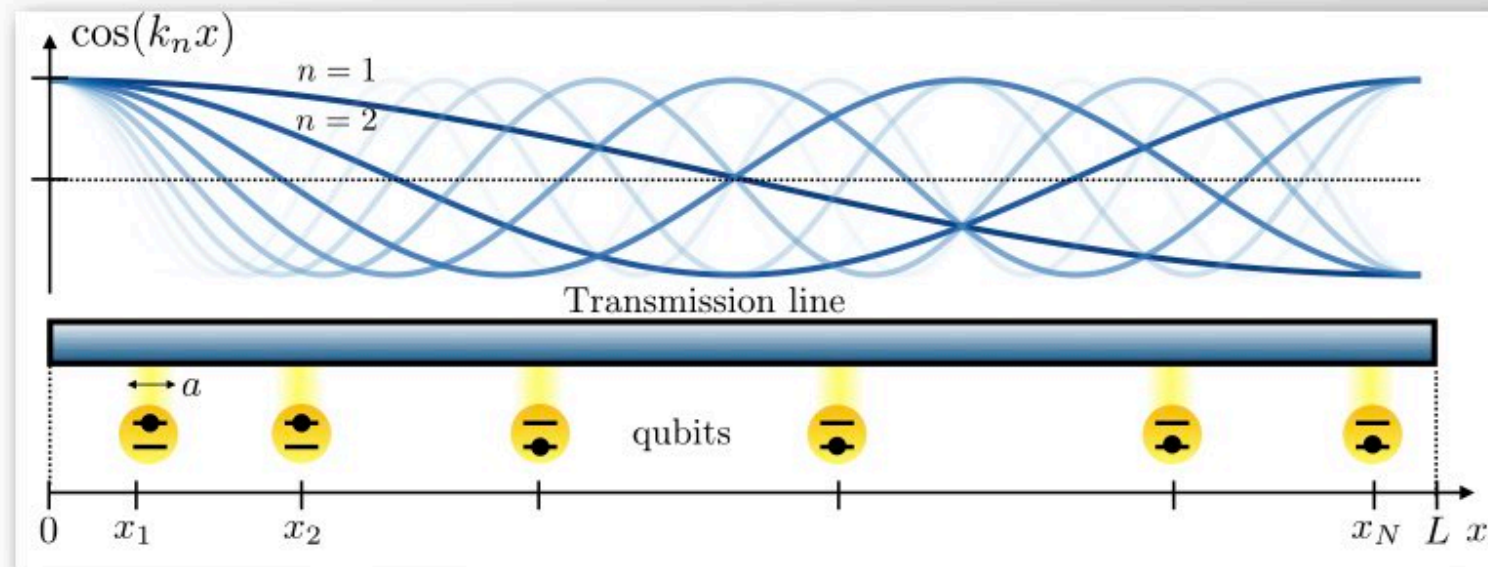
Matrix-Product State (25 Bosonic Modes)



Random noise in the transmission line



## Quantum Simulation of Spin Models



$$U_{\text{lab}}(t_p) = e^{-it_p \sum_i (\omega_i/2) \sigma_i^z} e^{-it_p \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z}.$$

$$J_{ij} = -2 \sum_n \frac{g_{i,n} g_{j,n}}{\omega_n}.$$

How to program an arbitrary interaction matrix?



Generate targeted and scalable time-evolution:

$$W = \exp(-i \sum_{i < j} w_{ij} \sigma_i^z \sigma_j^z)$$

$$w_{ij} = \sum_{q=1}^N w_q u_{i,q} u_{j,q}$$

Consider sequence of successive cycles:  $q = 1, \dots, \eta$

- Adjustable parameters:  $g_i \rightarrow g_i^{(q)}$
- Evolution at end of sequence:

$$U_\eta = \exp(-it_p \sum_{i < j} J_{i,j}^{(\eta)} \sigma_i^z \sigma_j^z)$$

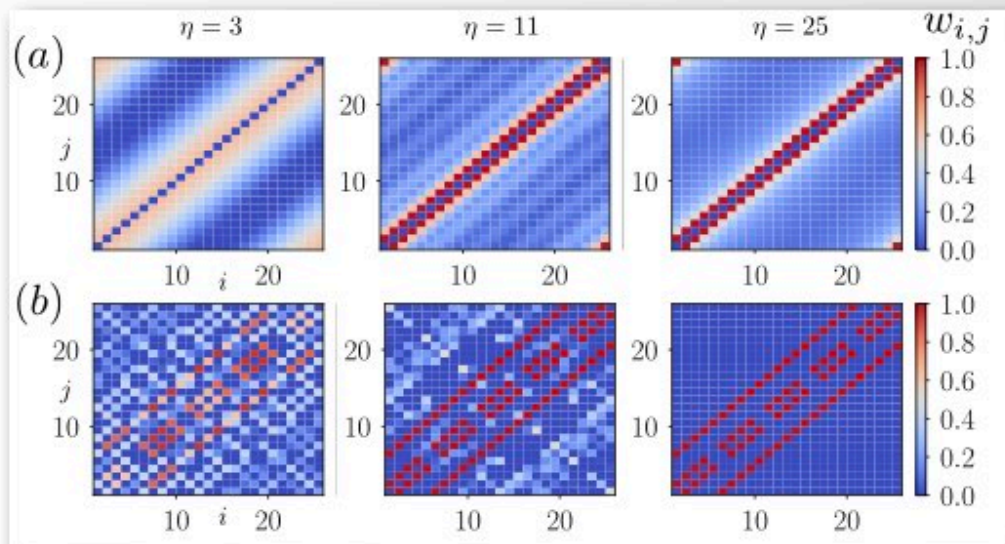
$$J_{i,j}^{(\eta)} = \sum_{q=1}^{\eta} \frac{g_i^{(q)} g_j^{(q)}}{\omega_1}$$

- Identify: This is our **recipe** for physical system.

$$g_i^{(q)} = \sqrt{w_q \omega_1 / t_p} u_{i,q}$$

- Engineer arbitrary spin-spin interactions within linear runtime:  $T = N t_p$

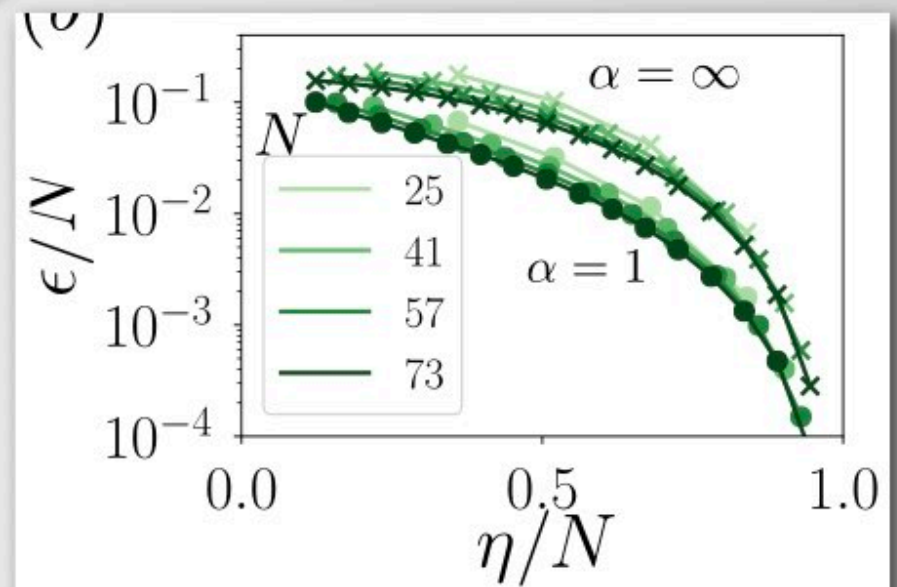
# Engineering of Spin Models



$$N = 25$$

Figure 3: *Engineering of spin models.* (a) Long range interactions  $w_{ij} = 1/|i - j|$  and periodic boundary conditions. (b) 2D nearest neighbor interactions with open boundary conditions. Here, the indices  $i$  correspond to 2D indices  $\mathbf{i} = (i_x, i_y)$  of a square of  $5 \times 5$  sites using the convention  $i = i_x + 5i_y$ .

- Can generate general Ising spin models in any spatial dimension and geometry.
- Observe progressive emergence of target matrix.



# Quantum Optimization (QAOA)

## QAOA Wavefunction

$$|\gamma, \beta\rangle = U_x(\beta_p)U_{zz}(\gamma_p) \cdots U_x(\beta_1)U_{zz}(\gamma_1)|s\rangle$$

$$U_x(\beta_m) = \exp[-i\beta_m \sum_i \sigma_i^x]$$

[Hot-gate protocol]

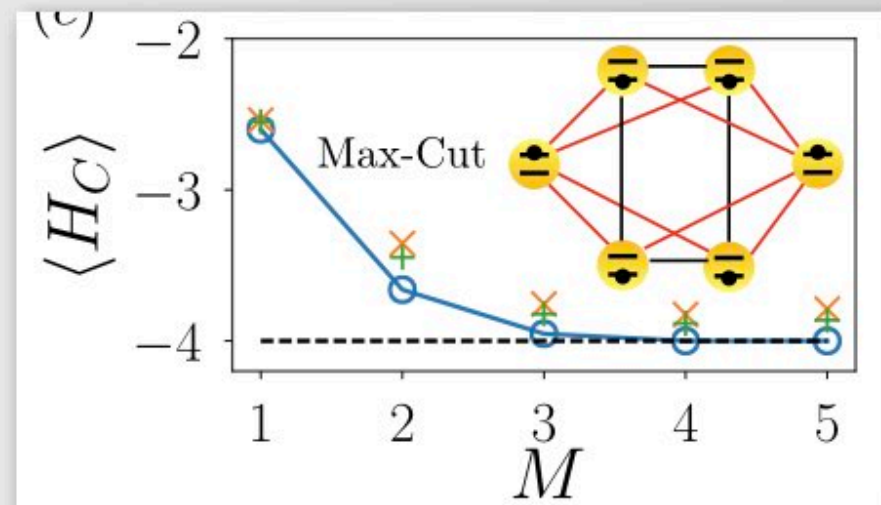
$$U_{zz}(\gamma_m) = \exp[-i\gamma_m H_C]$$

## Quantum Algorithm

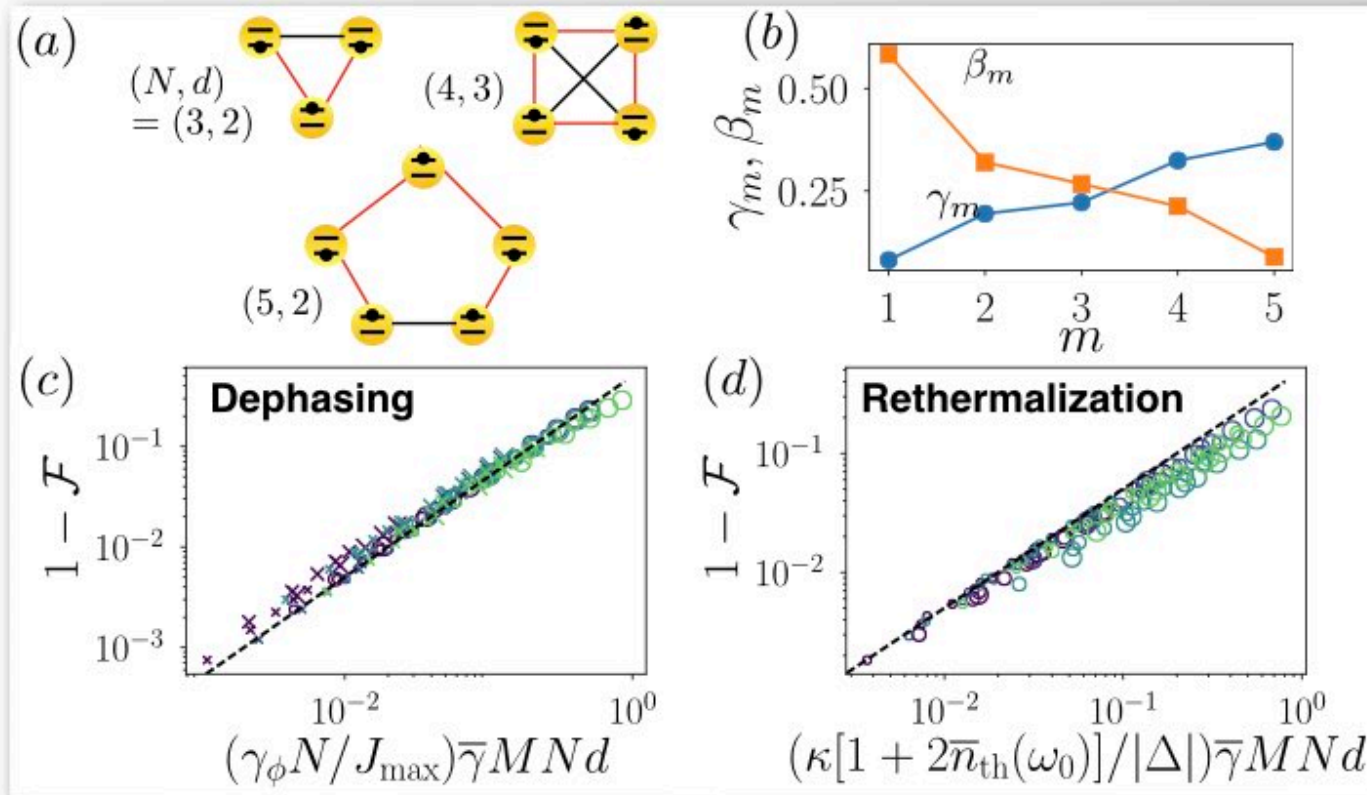
Implement QAOA wavefunction  
Measure Energy (spin configs)  
Get new Parameters  
Repeat until energy is conserved

Numerical simulations of our setup:

- Account for finite temperature.
- Account for decoherence.



# Robustness against Noise



Total QAOA error:

$$\xi \approx \bar{\gamma} d M N^{3/2} / \sqrt{C},$$

$$C = g^2 / (\gamma_\phi \kappa_{\text{eff}})$$

## **Model and exact Solution**

### **Illustrations**

Two-qubit gate

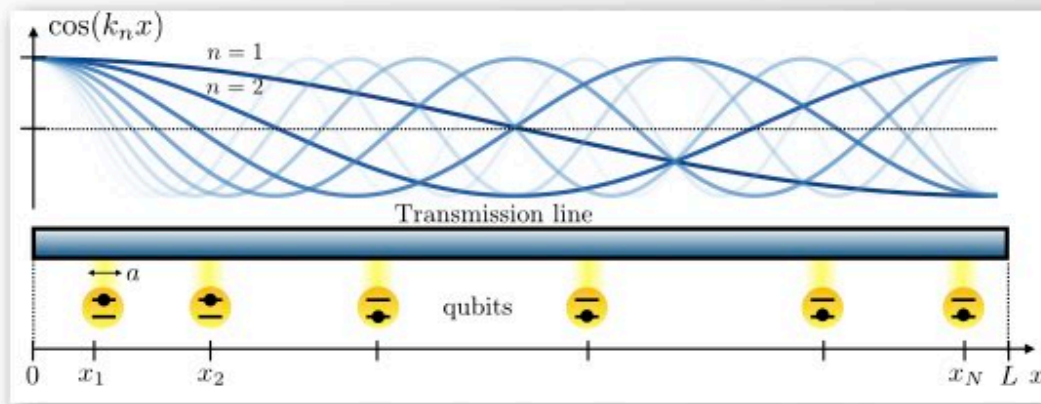
Quantum Simulation

Quantum optimization

## **Implementations with Superconducting qubits**



# Implementation



$$H = \sum_i \frac{\omega_i}{2} \sigma_i^z + \sum_{n=1}^{\infty} \omega_n a_n^\dagger a_n + \sum_{i,n} g_{i,n} \sigma_i^z (a_n + a_n^\dagger)$$

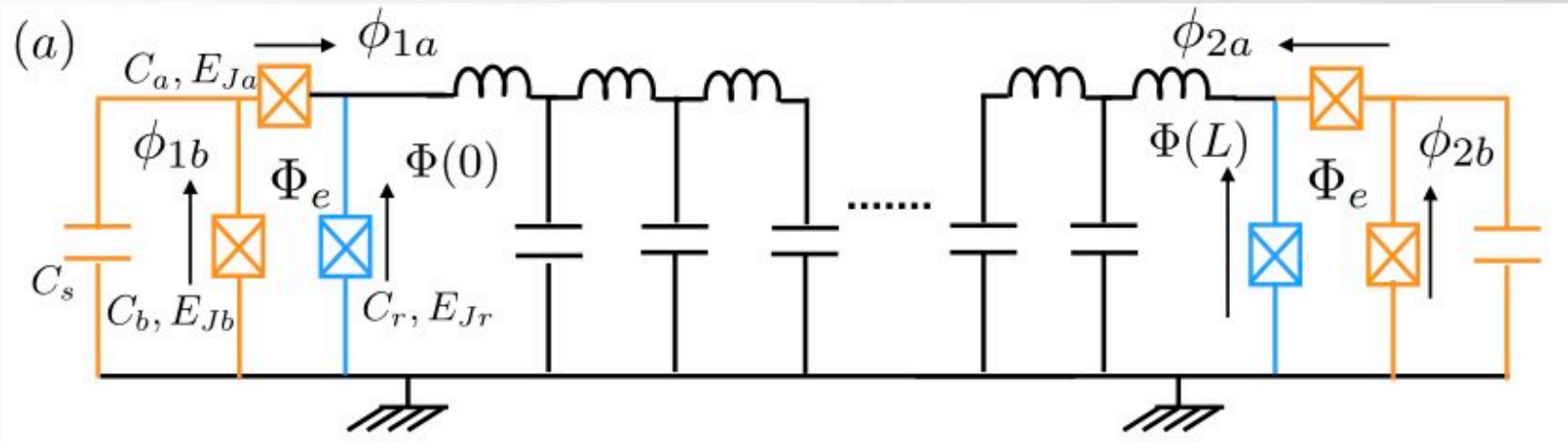
Longitudinal couplings

**First ideas in SC qubits:** Kerman-Nakamura-Blais (Phase Gates) - Blais (Readout)



**Our implementation is just a multi-qubit, multi-mode version**

# Implementation with two qubits



## Lagrangian

$$\mathcal{L} = \int_0^L dx \left( \frac{c\dot{\Phi}^2}{2} - \frac{(\partial_x \Phi)^2}{2\ell} \right)$$

TLine

$$+ \sum_{i=1,2} \left[ E_{Jr} \cos \left( \frac{\Phi(x_i)}{\phi_0} \right) + C_r \frac{\dot{\Phi}(x_i)^2}{2} \right]$$

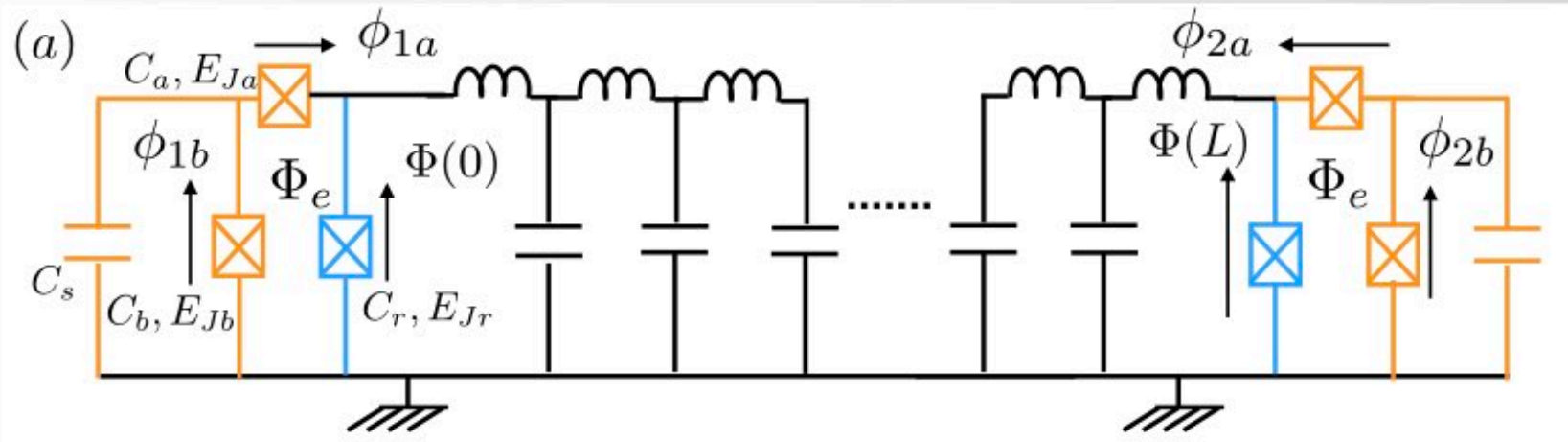
Coupling inductance

$$+ \frac{C_s + C_b}{2} \left( \frac{\dot{\Phi}(x_i)}{2} + \dot{\phi}_i \right)^2 + \frac{C_a}{2} \left( \frac{\dot{\Phi}(x_i)}{2} - \dot{\phi}_i \right)^2$$

Transmons

$$+ E_J \cos \left( \frac{\delta_i}{2\phi_0} \right) \cos \left( \frac{\phi_i}{\phi_0} \right) \Big], \quad (\text{S51})$$

# Implementation with two qubits

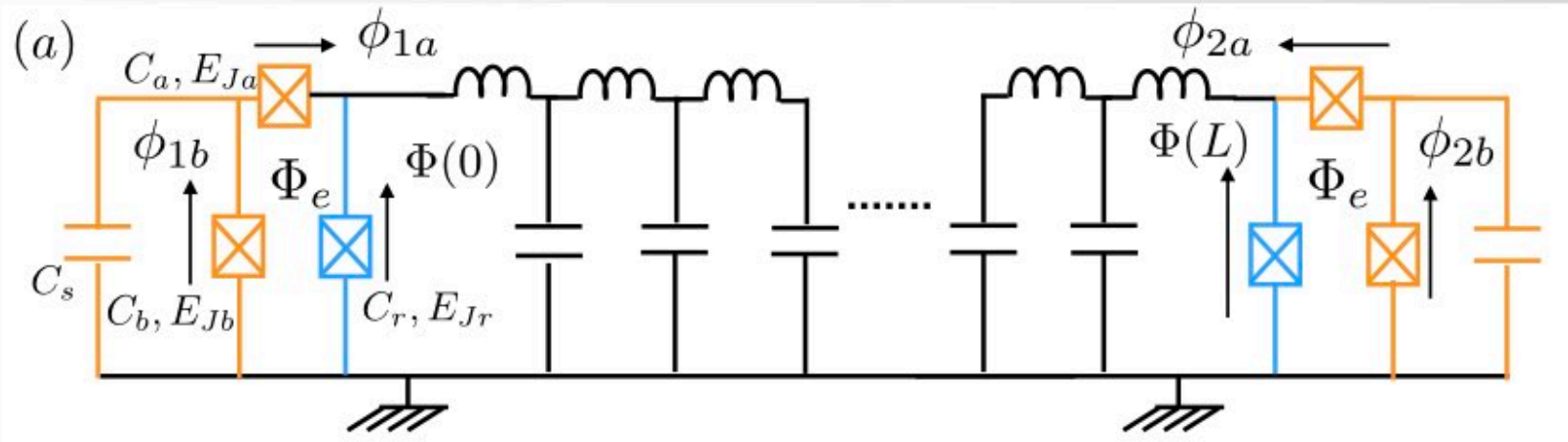


## Quantization

$$\Phi_n = \sqrt{\frac{\hbar}{2C\omega_n}}(a_n + a_n^\dagger), \quad q_n = \sqrt{\frac{\hbar C\omega_n}{2}}i(a_n^\dagger - a_n)$$

$$\phi_i = \sqrt{\frac{\hbar}{2C_T\omega_z}}(a_i + a_i^\dagger), \quad q_i = \sqrt{\frac{\hbar C_T\omega_z}{2}}i(a_i^\dagger - a_i),$$

# Implementation with two qubits



## Hamiltonian

**Driving:**

can be absorbed via displacement transformations

**Longitudinal coupling**

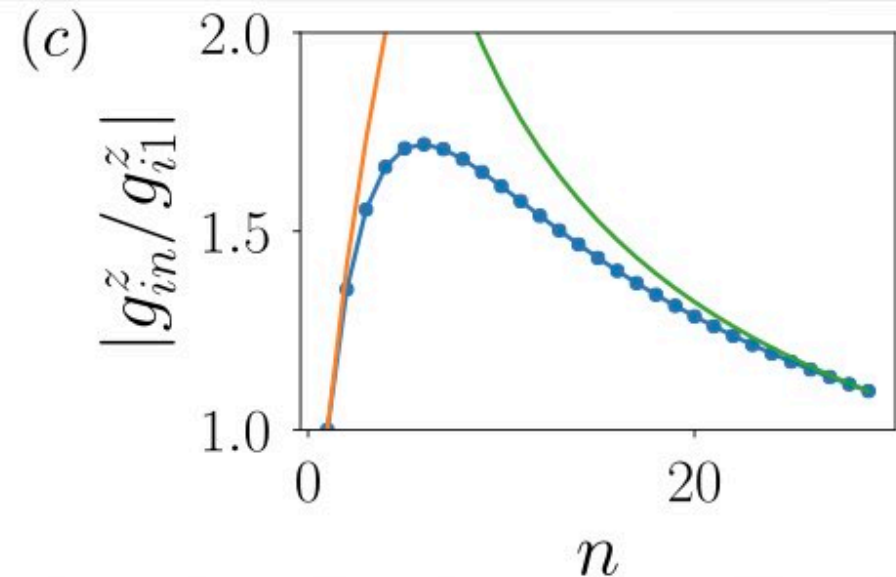
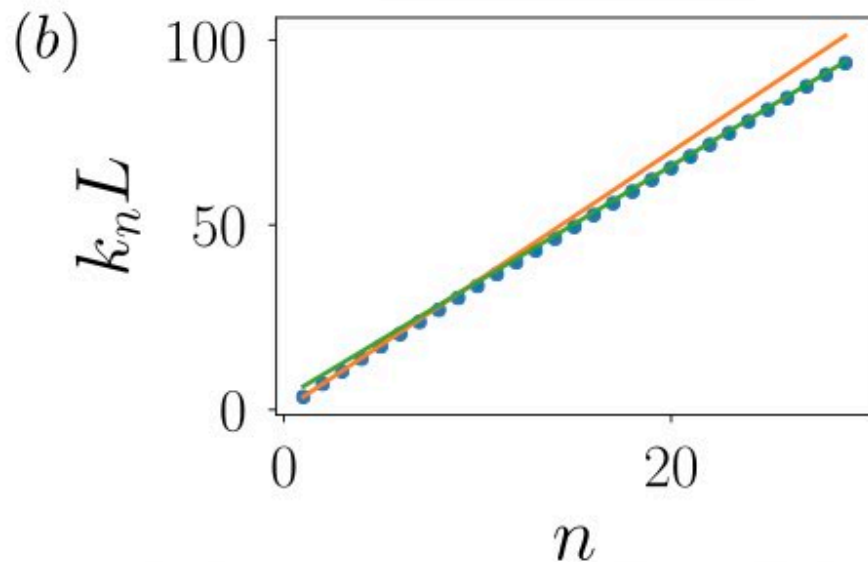
$$\begin{aligned}
 H_{\text{int}} = & \hbar \sum_{i,n} \Omega_{i,n} (a_n^\dagger + a_n) + \hbar \sum_{i,n} g_{in}^z \sigma_i^z (a_n^\dagger + a_n), \\
 & + \hbar \sum_{i,n} g_{in}^y (a_i^\dagger - a_i)(a_n^\dagger - a_n), \tag{S62}
 \end{aligned}$$

**Transverse coupling:**

can be eliminated with condition on the capacitance

# Implementation with two qubits

## Numerical results



👍 Quasi-Linear Dispersion Relation    👍 Coupling strengths with cutoff

Typical Energy Scale (60 MHz)



Thank you!!!!



**M.J.A. Schuetz**



M.D. Lukin



L.M.K. Vandersypen



G. Kirchmair



P. Zoller



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